



NTNU
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Computation of the Hessian for Full Waveform Inversion

Vegard Stenhjem Hagen

Iterative methods¹

- Searching for a model \mathbf{m} that describes the earth.
- Solve the wave equation

$$\frac{\partial^2}{\partial t^2} \mathbf{u} - \mathbf{c}^2 \nabla^2 \mathbf{u} = \mathbf{f}, \quad \mathbf{m} = [\mathbf{c}],$$

where \mathbf{u} is displacement and the model \mathbf{m} only consists of velocities \mathbf{c} , \mathbf{f} is a source function.

- Compare with true data \mathbf{u}_0 using a misfit function

$$\Psi(\mathbf{u}(\mathbf{m}, \mathbf{x}_r), \mathbf{u}_0) = \frac{1}{2} (\mathbf{u}_0(\mathbf{x}_r) - \mathbf{u}(\mathbf{m}, \mathbf{x}_r))^T (\mathbf{u}_0(\mathbf{x}_r) - \mathbf{u}(\mathbf{m}, \mathbf{x}_r)).$$

- Iterative approach. Find a model update $\delta \mathbf{m}_k$ that decreases the misfit

$$\Psi(\mathbf{m}_{k+1} = \mathbf{m}_k + \delta \mathbf{m}_k) < \Psi(\mathbf{m}_k).$$

¹Tarantola 1984; Mora 1987; Fichtner et al. 2006; Fichtner 2011.

Iterative methods

- To do this we can calculate the gradient of the misfit

$$\mathbf{J}(\mathbf{m} + \delta\mathbf{m}) = \nabla_m \Psi(\mathbf{m} + \delta\mathbf{m}).$$

- By linearising around the Jacobian we get

$$\mathbf{J}(\mathbf{m} + \delta\mathbf{m}) \simeq \mathbf{J}(\mathbf{m}) + \nabla_m \mathbf{J}(\mathbf{m}) \delta\mathbf{m} = \mathbf{0}.$$

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- The Hessian is then given by

$$\mathbb{H}(\mathbf{m}) = \nabla_m \mathbf{J}(\mathbf{m}) = \nabla_m \nabla_m \Psi(\mathbf{m}).$$

Newton method²

- By solving

$$\mathbb{H}(\mathbf{m})\delta\mathbf{m} = -\mathbf{J}(\mathbf{m})$$

for $\delta\mathbf{m}$ we find the next model update.

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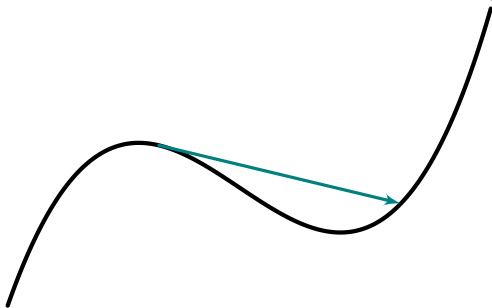
- A common approximation is

$$\delta\mathbf{m} \simeq \alpha\mathbf{J},$$

and a line search for the optimal $\alpha \in \mathbf{R}$.

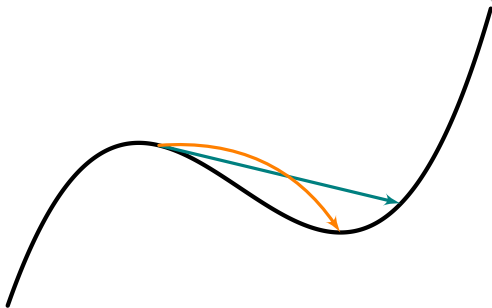
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Advantages of using the full Hessian



Linear approximation

Advantages of using the full Hessian



Adjust for slope

Advantages of using the full Hessian³

- Fewer iterations are expected.
- Possible to perform resolution analysis.
- Parameter cross-talk analysis.
- Incorporate second order effects.
- Extra transmission-like information.

³Pratt et al. 1998; Fichtner and Trampert 2011b; Trampert et al. 2013; Biondi et al. 2015.

Fréchet derivative

- Use the adjoint approach by Tarantola 1984 and further work in Fichtner and Trampert 2011a.
- The Fréchet derivative is defined as

$$\nabla_m \Psi(\mathbf{m}) \delta \mathbf{m} = \lim_{\nu \rightarrow 0} \frac{1}{\nu} [\Psi(\mathbf{m} + \nu \delta \mathbf{m}) - \Psi(\mathbf{m})],$$

i.e. the derivative of Ψ with respect to \mathbf{m} in the $\delta \mathbf{m}$ direction.

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- The Jacobian acting on a model perturbation $\delta \mathbf{m}$ we can be written as

$$\mathbf{J}(\mathbf{m}) \delta \mathbf{m} = \nabla_m \Psi(\mathbf{m}) \delta \mathbf{m}.$$

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- The Jacobian acting on a model perturbation $\delta \mathbf{m}$ we can be written as

$$\mathbf{J}(\mathbf{m}) \delta \mathbf{m} = \nabla_m \Psi(\mathbf{m}) \delta \mathbf{m}.$$

- For the Hessian \mathbb{H} we can calculate its action on a model perturbation $\delta \mathbf{m}$ as

$$\mathbb{H}(\mathbf{m}) \delta \mathbf{m} = \nabla_m \nabla_m \Psi(\mathbf{m}) \delta \mathbf{m}.$$

Fréchet kernel

- The Fréchet kernel is defined as the volumetric densities of the Fréchet derivative

$$\mathbf{F} = \frac{d}{dV} \nabla_m \Psi$$

giving us the relation

$$\mathbf{J} \delta \mathbf{m} = \nabla_m \Psi \delta \mathbf{m} = \int_G \mathbf{F} \delta \mathbf{m} d^3 \mathbf{x}, \quad G \in \mathbb{R}^3.$$

- By studying \mathbf{F} we can design more efficient inversion schemes and interpret the results in a meaningful way.

Fréchet kernel

- Introducing the adjoint wavefield \mathbf{u}^\dagger defined as the solution to

$$\frac{\partial^2}{\partial t^2} \mathbf{u}^\dagger - \mathbf{c}^2 \nabla^2 \mathbf{u}^\dagger = -(\mathbf{u}_0(\mathbf{x}_r) - \mathbf{u}(\mathbf{x}_r)),$$

which is referred to as backpropagating the residuals.

- The Fréchet kernel can now be calculated as

$$\mathbf{F}(\mathbf{u}^\dagger, \mathbf{u}) = 2\mathbf{c} \int_T \nabla \mathbf{u}^\dagger \nabla \mathbf{u} dt,$$

i.e. by cross correlating the divergence of the forward and adjoint displacement fields and scaling by the background velocity.

Perturbed fields

- In order to more easily calculate the Hessian we need to introduce two more fields.
- The perturbed forward field

$$\begin{aligned}\delta \mathbf{u} &= \nabla_m \mathbf{u} \delta \mathbf{m} \\ &= \lim_{\nu \rightarrow 0} \frac{1}{\nu} [\mathbf{u}^\dagger(\mathbf{m} + \nu \delta \mathbf{m}) - \mathbf{u}^\dagger(\mathbf{m})]\end{aligned}$$

- The perturbed adjoint field

$$\begin{aligned}\delta \mathbf{u}^\dagger &= \nabla_m \mathbf{u}^\dagger \delta \mathbf{m} \\ &= \lim_{\nu \rightarrow 0} \frac{1}{\nu} [\mathbf{u}(\mathbf{m} + \nu \delta \mathbf{m}) - \mathbf{u}(\mathbf{m})]\end{aligned}$$

Hessian kernels

- The Hessian kernel \mathbf{H} can be broken down into three parts

$$\mathbf{H} = \mathbf{H}^{\mathbf{u}^\dagger, \delta \mathbf{u}} + \mathbf{H}^{\delta \mathbf{u}^\dagger, \mathbf{u}} + \mathbf{H}^{\mathbf{u}^\dagger, \mathbf{u}}.$$

- Written out these are

$$\mathbf{H}^{\mathbf{u}^\dagger, \delta \mathbf{u}} = 2\mathbf{c} \int_T \nabla(\mathbf{u}^\dagger) \cdot \nabla(\delta \mathbf{u}) dt$$

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$$\mathbf{H}^{\mathbf{u}^\dagger, \delta \mathbf{u}} = 2\mathbf{c} \int_T \nabla(\mathbf{u}^\dagger) \cdot \nabla(\delta \mathbf{u}) dt = \mathbf{F}(\mathbf{u}^\dagger, \delta \mathbf{u}),$$

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$$\mathbf{H}^{\mathbf{u}^\dagger, \mathbf{u}} = 2\delta \mathbf{c} \int_T \nabla(\mathbf{u}^\dagger) \cdot \nabla(\mathbf{u}) dt = \frac{\delta \mathbf{c}}{\mathbf{c}} \mathbf{F}(\mathbf{u}^\dagger, \mathbf{u}).$$

Fields

— We now have four fields we need to calculate.

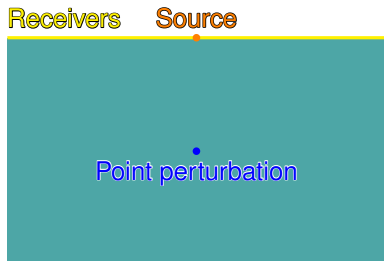
\mathbf{u} — Forward field.

\mathbf{u}^\dagger — Adjoint field.

$\delta\mathbf{u}$ — Perturbed forward field.

$\delta\mathbf{u}^\dagger$ — Perturbed adjoint field.

Model



- Acoustic 2-D
- Background 2.0 km/s
- Perturbation 2.2 km/s
- Width 1500 m, Depth 750 m
- 3.75 m \times 3.75 m grid cells
- 650 ms in 0.1 ms time steps
- 20 Hz Ricker wavelet

Forward field

Adjoint field

Perturbed forward field

Perturbed adjoint field

Fields

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Reflection - Fréchet kernel — $F(u^\dagger, u)$

Divergence Correlation

Background field

Adjoint field

Reflection — $\mathbf{H}^{\mathbf{u}^\dagger, \delta \mathbf{u}} = \mathbf{F}(\mathbf{u}^\dagger, \delta \mathbf{u})$

Divergence Correlation

Perturbed forward field

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Divergence Correlation

Background field

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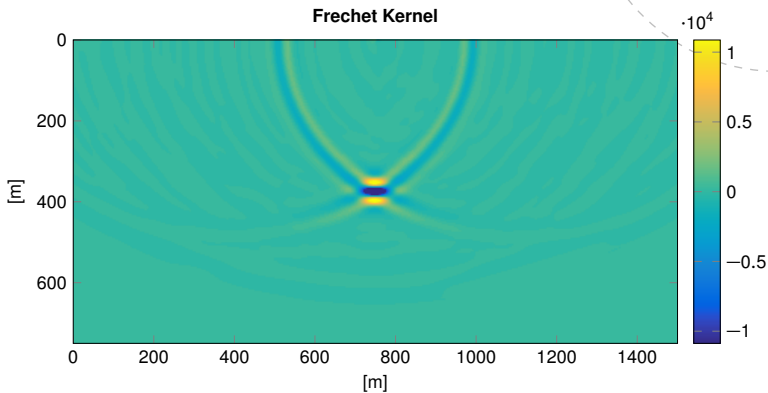
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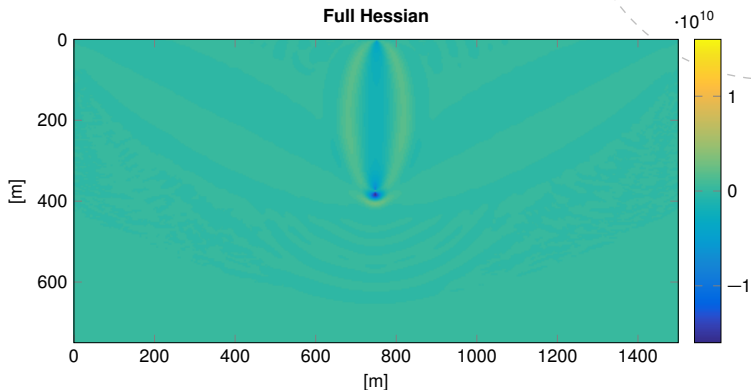
Background field

Adjoint field

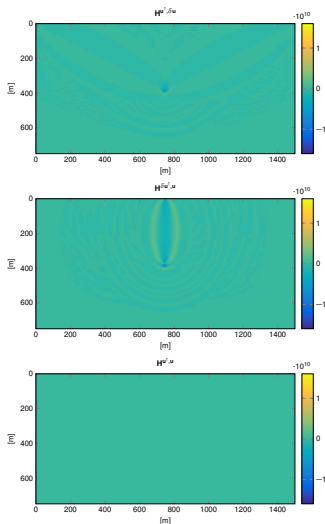
Reflection - Fréchet kernel — $F(u^\dagger, u)$



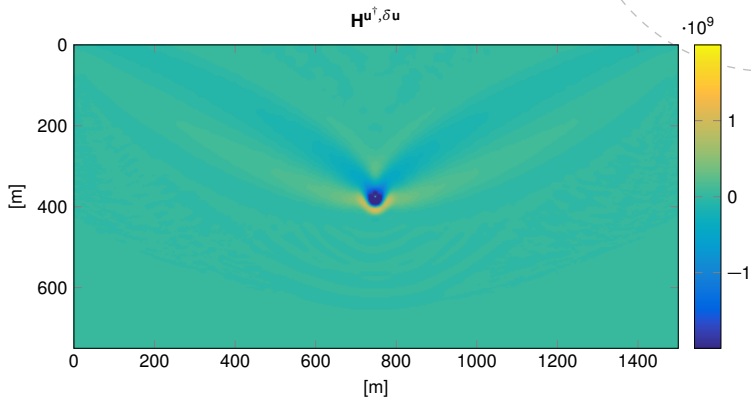
Reflection - Full Hessian Kernel — H



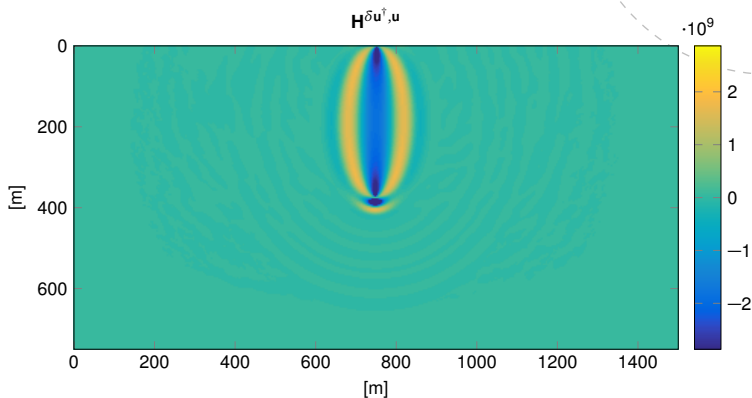
Reflection



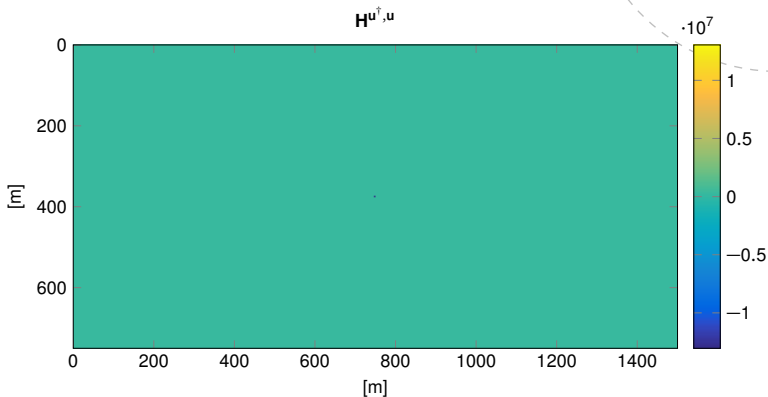
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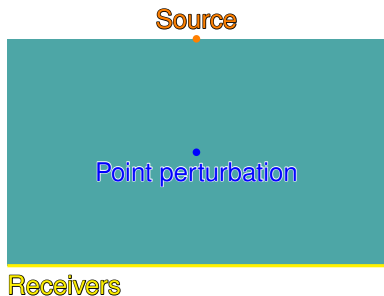
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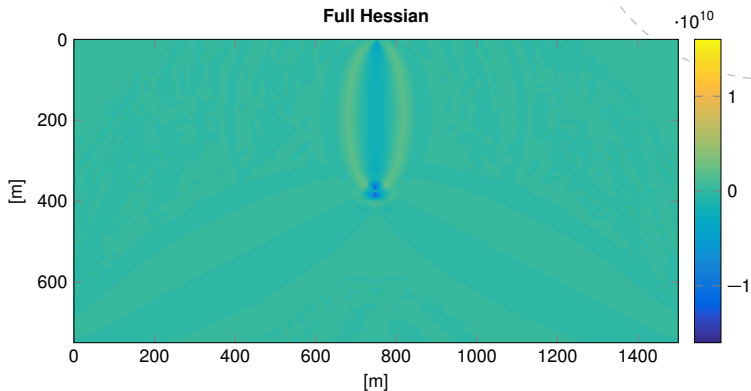


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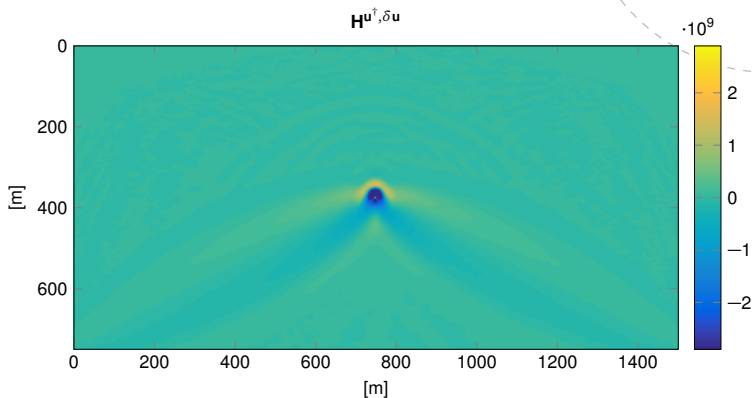


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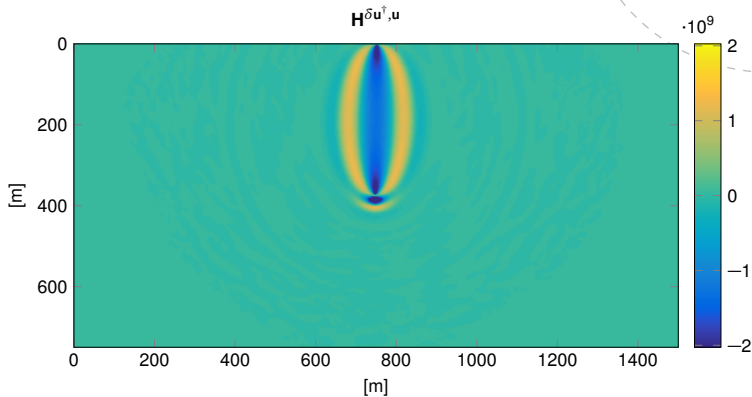
Transmission - Full Hessian Kernel — H



Transmission — $\mathbf{H}^{\mathbf{u}^\dagger, \delta \mathbf{u}} = \mathbf{F}(\mathbf{u}^\dagger, \delta \mathbf{u})$



Transmission — $\mathbf{H}^{\delta \mathbf{u}^\dagger, \mathbf{u}} = \mathbf{F}(\delta \mathbf{u}^\dagger, \mathbf{u})$



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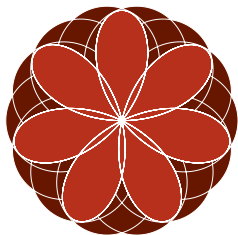
Future work

- Implement in 3-D for Madagascar.
- Calculate the Hessian for an elastic medium.
- Write a full Newton inversion algorithm.
- Investigate cross-talk in different parametrisations.

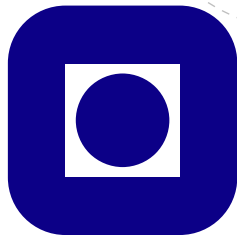
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Acknowledgements



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