

Static and Dynamic Moduli

ROSE

Rock Physics and Geomechanics

Course 2012

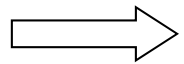
Erling Tjørr

What do we mean with "static" and "dynamic" moduli?

Elastic wave velocities

$$V_p = \sqrt{\frac{\lambda_e + 2G_e}{\rho}}$$

$$V_s = \sqrt{\frac{G_e}{\rho}}$$



$$G_e = \rho V_s^2$$

$$\lambda_e = \rho V_p^2 - 2\rho V_s^2$$

Dynamic moduli

Stress and strain measured
in a rock mechanical test

$$E = \frac{\Delta \sigma_z}{\Delta \varepsilon_z}$$

Static moduli

In general:

$$E_{stat} \neq E_{dyn}$$

$$E_{stat} < E_{dyn}$$

Note: $E_{stat} : E_{dyn}$ is not a constant ratio – it changes with stress!

In saturated rocks,

- Static deformation is often drained
- Dynamic moduli are always undrained

Occasionally used definition:

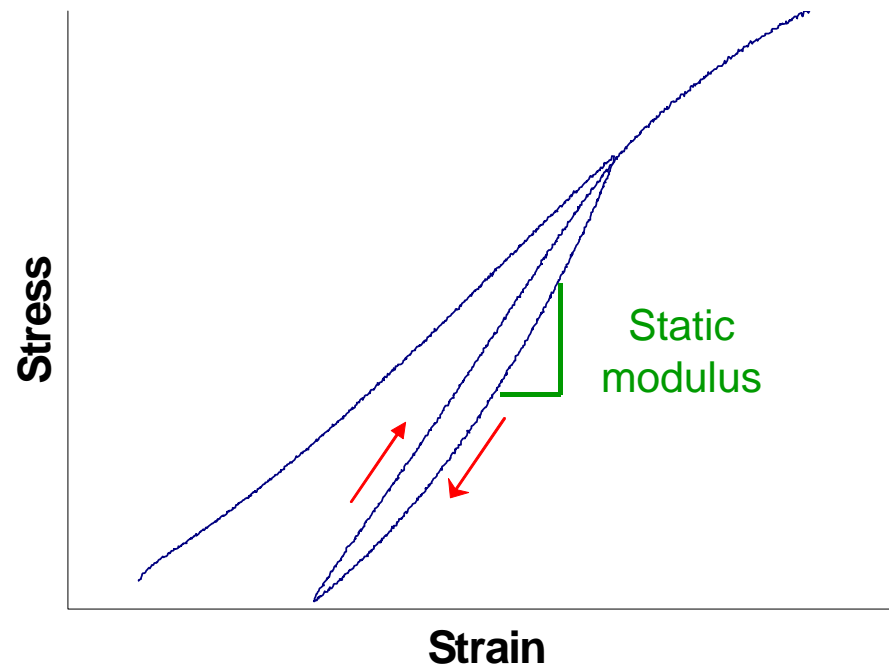
"Static modulus = drained modulus"

"Dynamic modulus" = undrained modulus"

Not recommended!

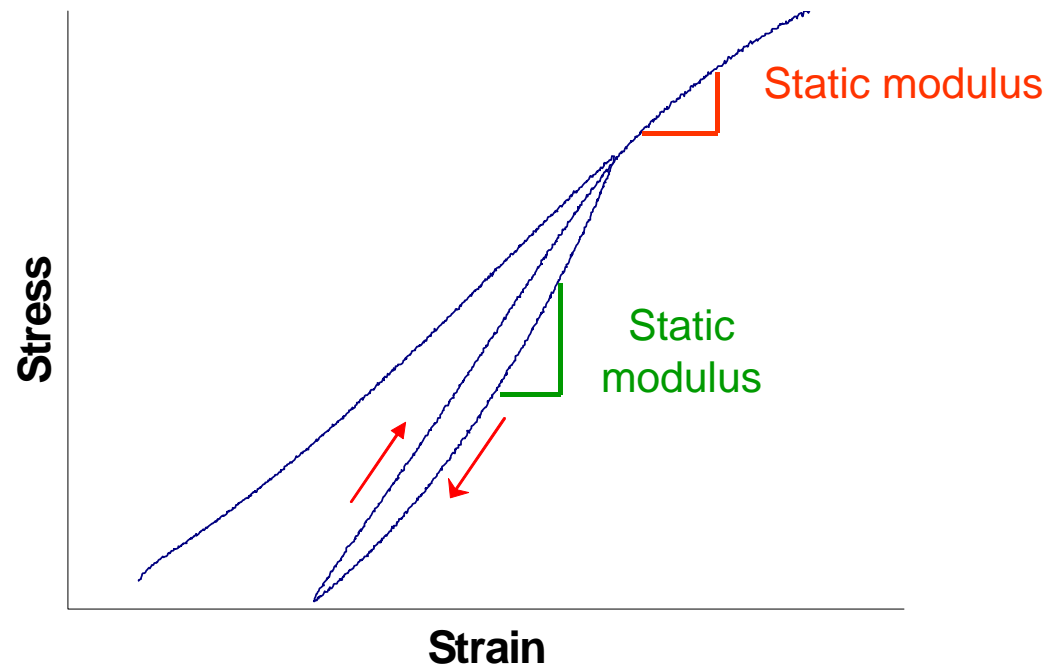
Alternative definition, also used:

"Static modulus" = slope of stress-strain curve measured during unloading

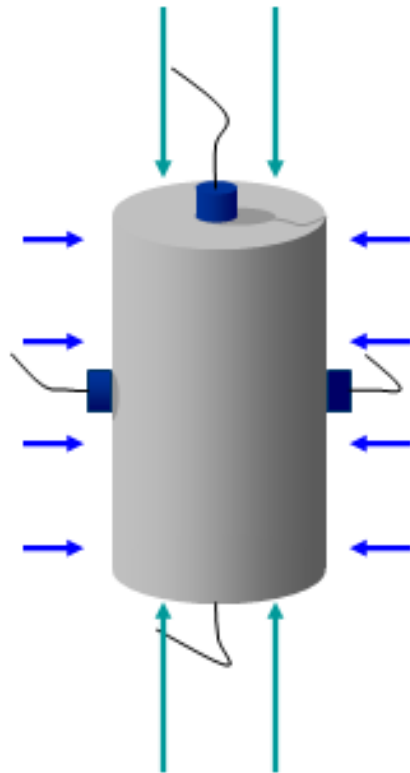


Our definition:

"Static modulus" = slope of stress-strain curve



Laboratory tests:



Standard triaxial set-up + acoustics

Measurements:

Stress

Strain

Acoustic wave velocities

Enables simultaneous measurements of

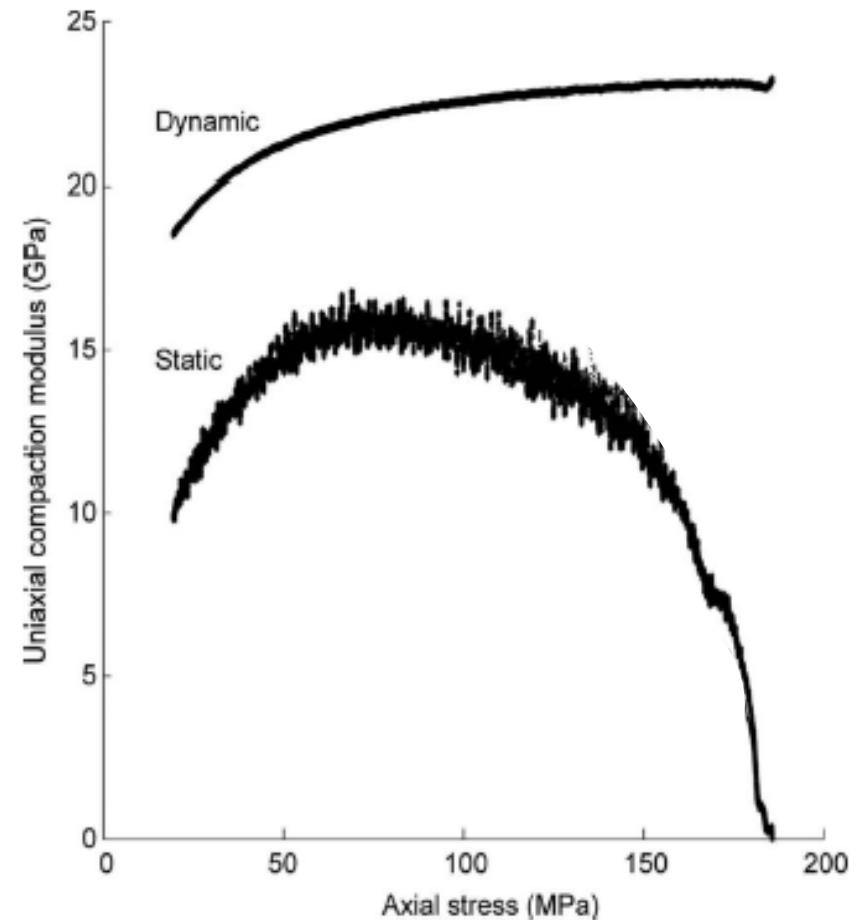
- static moduli (slope of stress-strain curve)
- dynamic moduli (density \times velocity²)

Static and dynamic moduli of soft rocks are different.

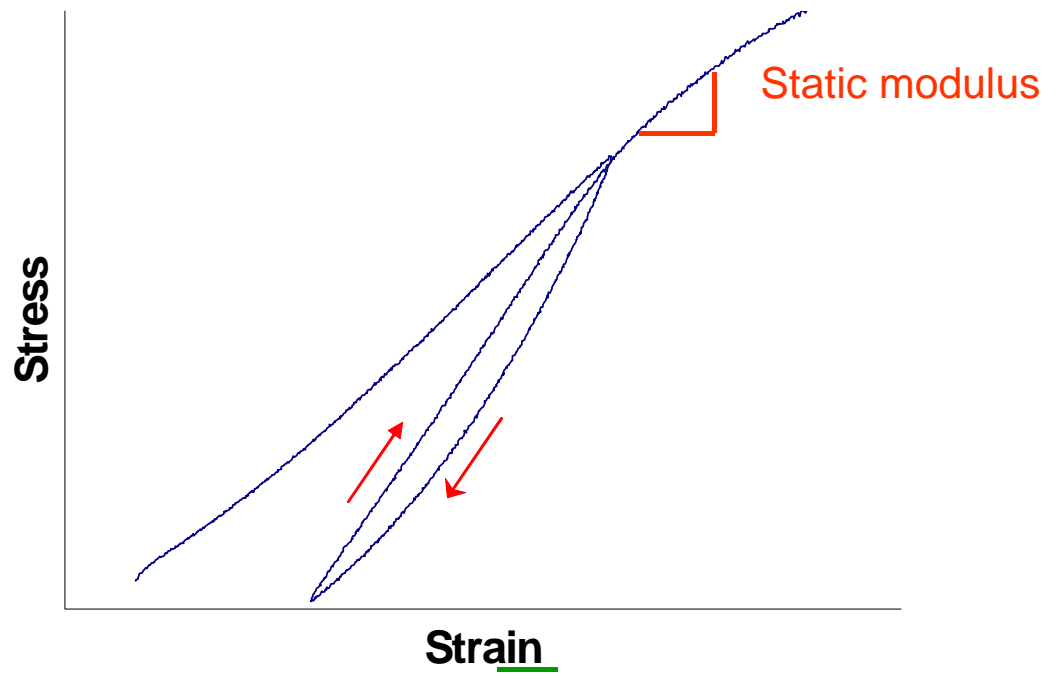
The difference changes along the stress path.

Potential causes for the difference between static and dynamic moduli:

- Strain rate
- Length of stress path
- Stress history
- Rock volume involved
- Drainage conditions
- Anisotropy



First: consider the static modulus measured during **initial loading**



Fjær (1999):

We introduce a **parameter P** , defined as:

$$P = \frac{\Delta\varepsilon_v - \Delta\varepsilon_{v,e}}{3\Delta\sigma}$$

P is a measure of the inelastic part of the deformation caused by a compressive hydrostatic stress increment.

$\Delta\varepsilon_v$ - total volumetric strain

$$\Delta\varepsilon_{v,e} \equiv \frac{\Delta\sigma}{K_e} \text{ - elastic strain}$$

⇒

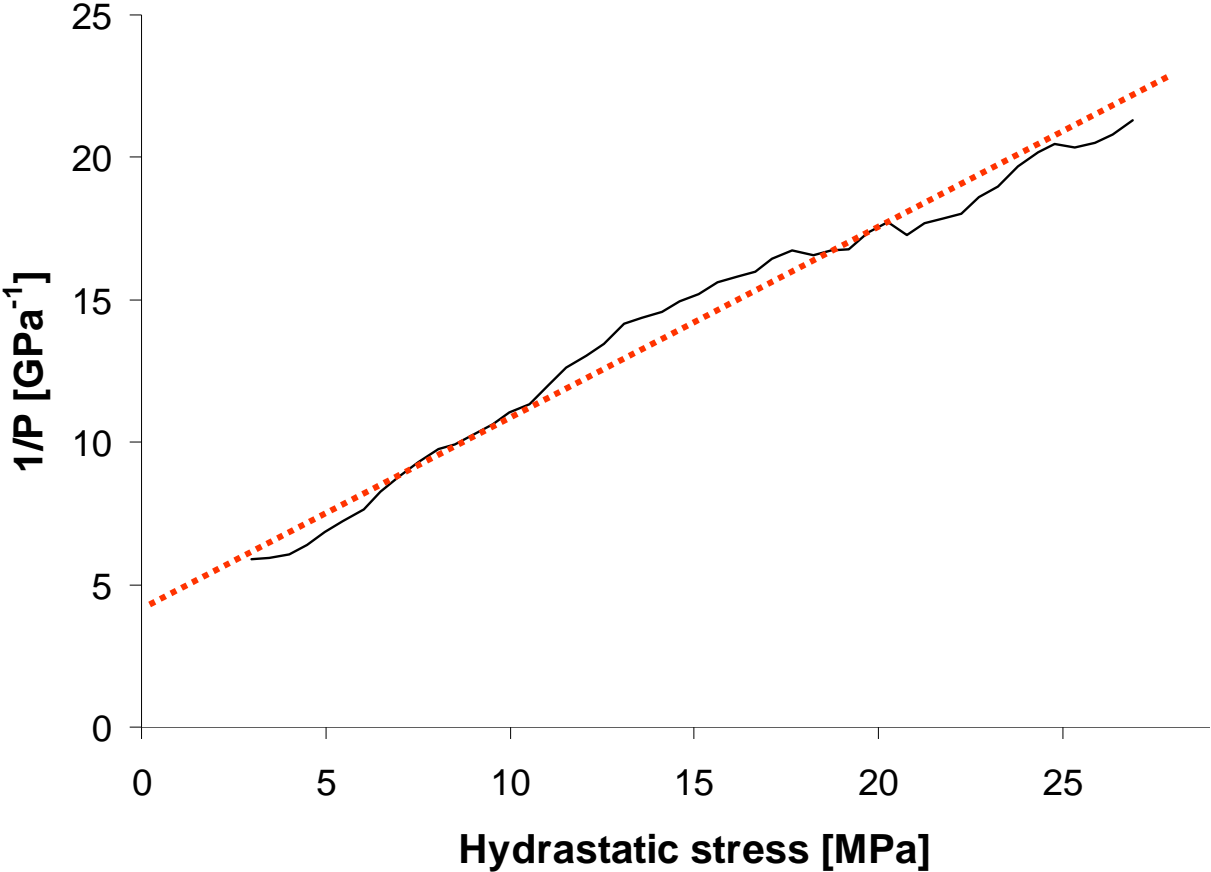
$$K = \frac{K_e}{1 + 3PK_e}$$

K = Static bulk modulus

K_e = Dynamic bulk modulus

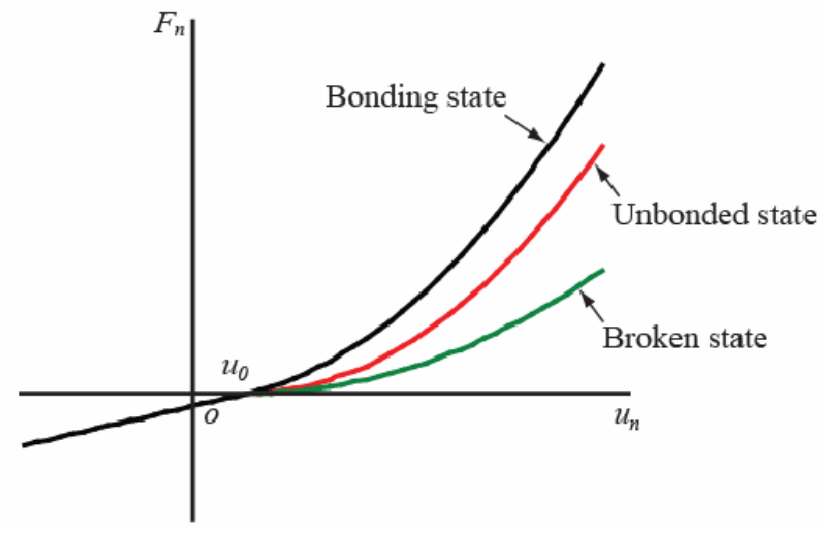
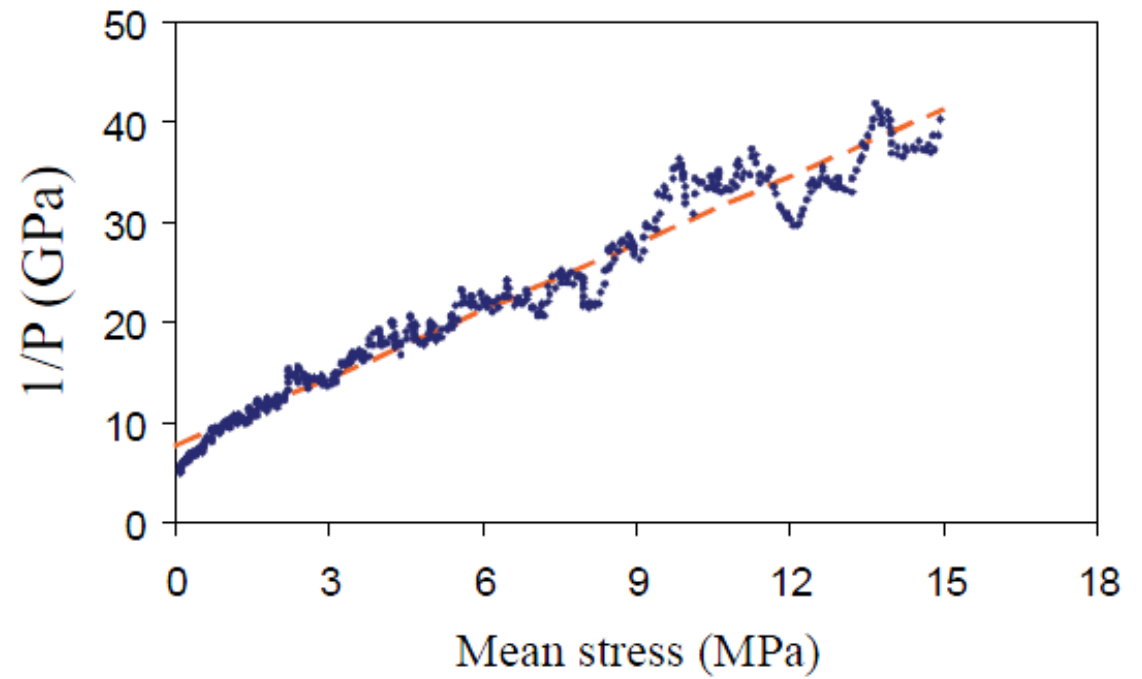
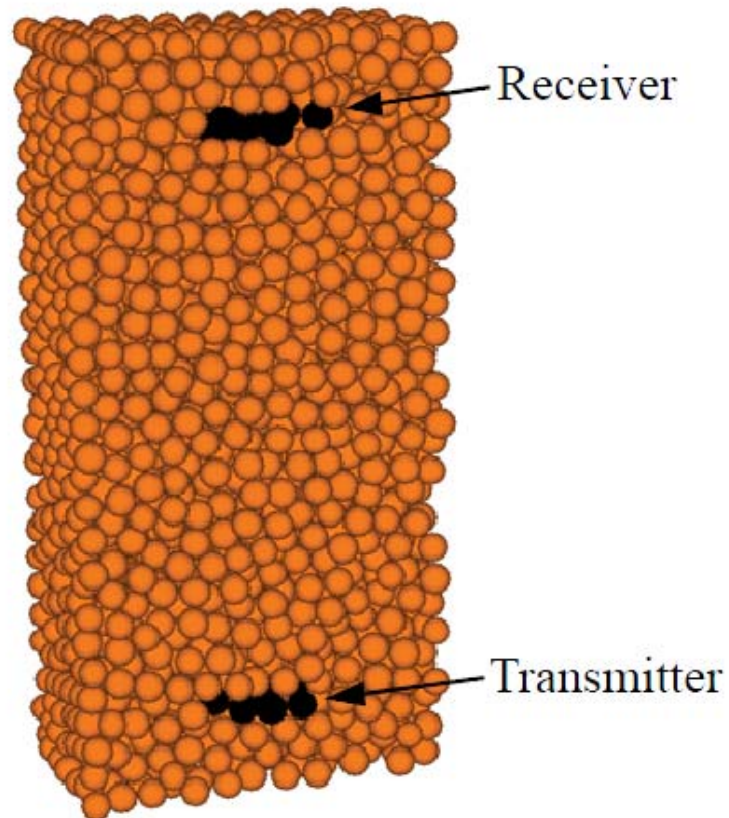
Observations

Hydrostatic test



$$P = \frac{\epsilon_g}{\sigma + T}$$

PFC^{3D}-simulation
Li & Fjær, 2008



Fjær (1999):

We introduce a **parameter F** , defined as:

$$F = \frac{\Delta \varepsilon_z - \Delta \varepsilon_{z,e} - \Delta \varepsilon_{z,p}}{\Delta \varepsilon_z}$$

F is a measure of the inelastic part of the deformation caused by a shear stress increment.

$\Delta \varepsilon_z$ - total axial strain

$$\Delta \varepsilon_{z,e} \equiv \frac{\Delta \sigma}{E_e} \text{ - elastic strain}$$

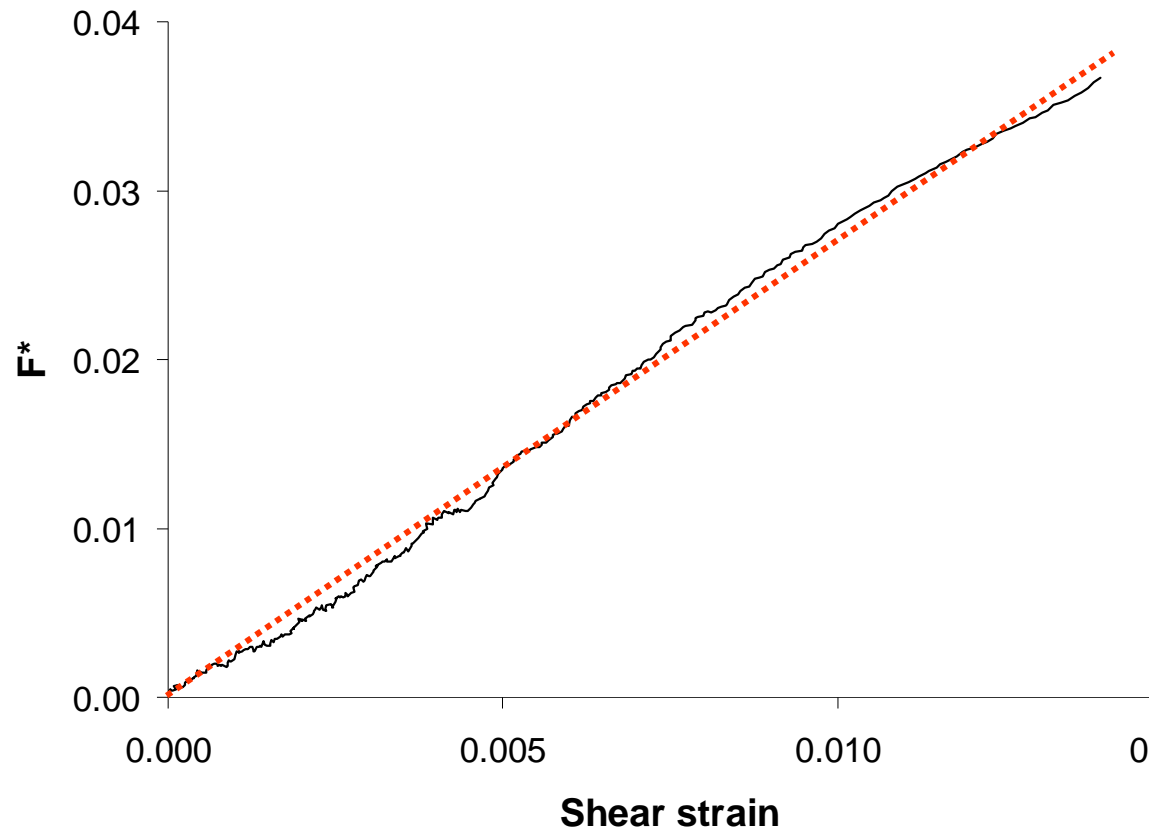
$$\Delta \varepsilon_{z,p} = P_z \Delta \sigma_z$$

⇒

$$E = \frac{E_e}{1 + P_z E_e} (1 - F)$$

E = Static Young's modulus
 E_e = Dynamic Young's modulus

$$F^* = F \sqrt{\sigma_z + \sigma_r + S}$$



$$F = A \frac{\epsilon_z - \epsilon_r - \epsilon_o}{\sqrt{\sigma_z + \sigma_r + S}}$$

Discussion: the F - parameter

$$E = \frac{E_e}{1 + P_z E_e} (1 - F)$$

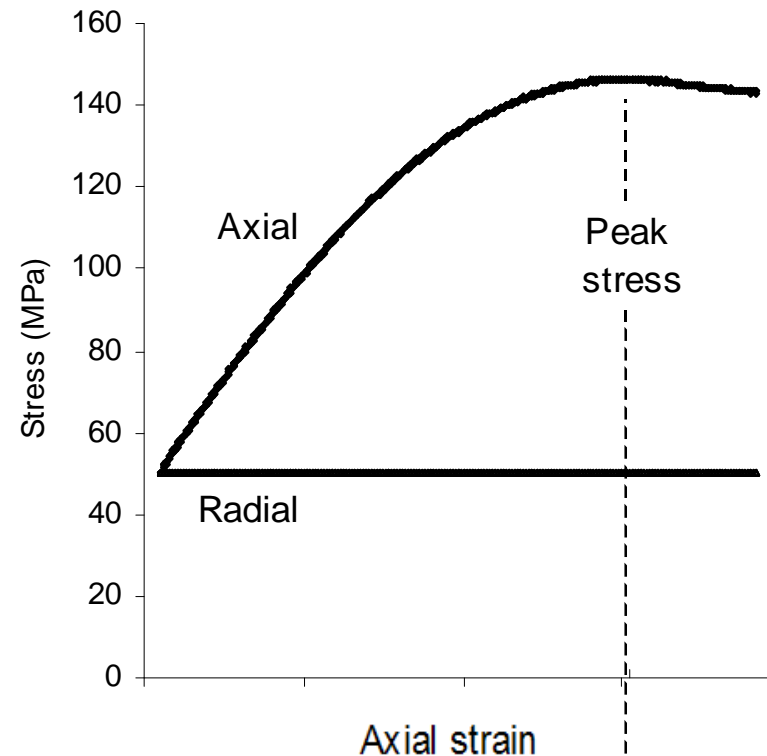
Note:

Since $E \propto (1 - F)$

\Rightarrow when $F = 1$ then $E = 0$

\Rightarrow peak stress

$F = 1 \leftrightarrow$ rock strength



We have a set of equations.....

These represent a constitutive model for the rock

We may use it to predict rock behavior, and thereby derive mechanical properties for the rock

$$K = \frac{K_e}{1 + 3PK_e}$$

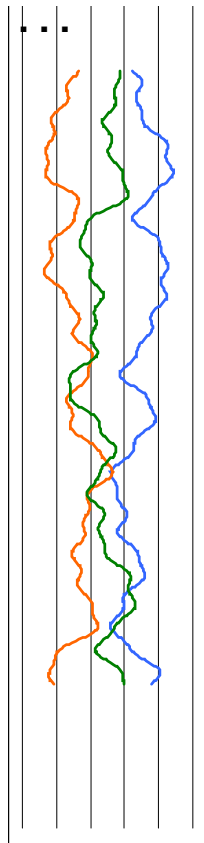
$$E = \frac{E_e}{1 + P_z E_e} (1 - F)$$

$$P = \frac{\varepsilon_g}{\sigma + T}$$

$$F = A \frac{\varepsilon_z - \varepsilon_r - \varepsilon_o}{\sqrt{\sigma_z + \sigma_r + S}}$$

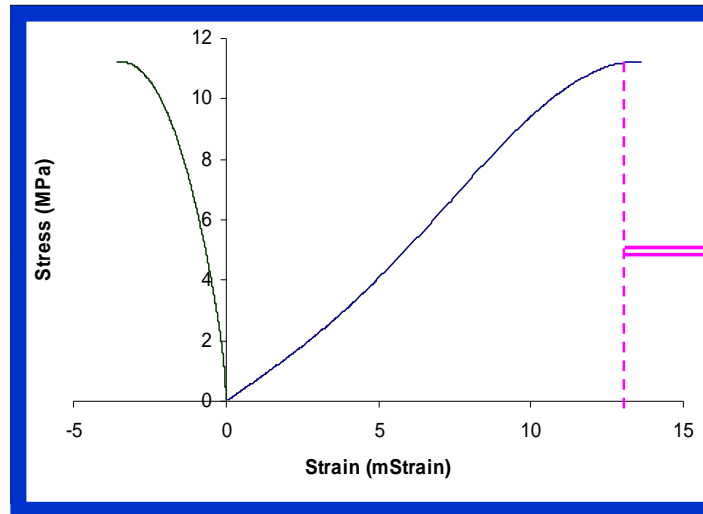
Application for logging purposes

Porosity,
Density, Sonic, .



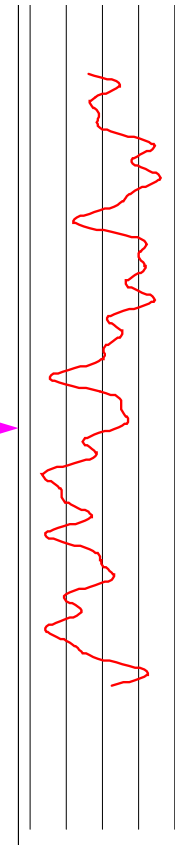
Provides calibration for the model

Constitutive model



Provides strength and stiffness

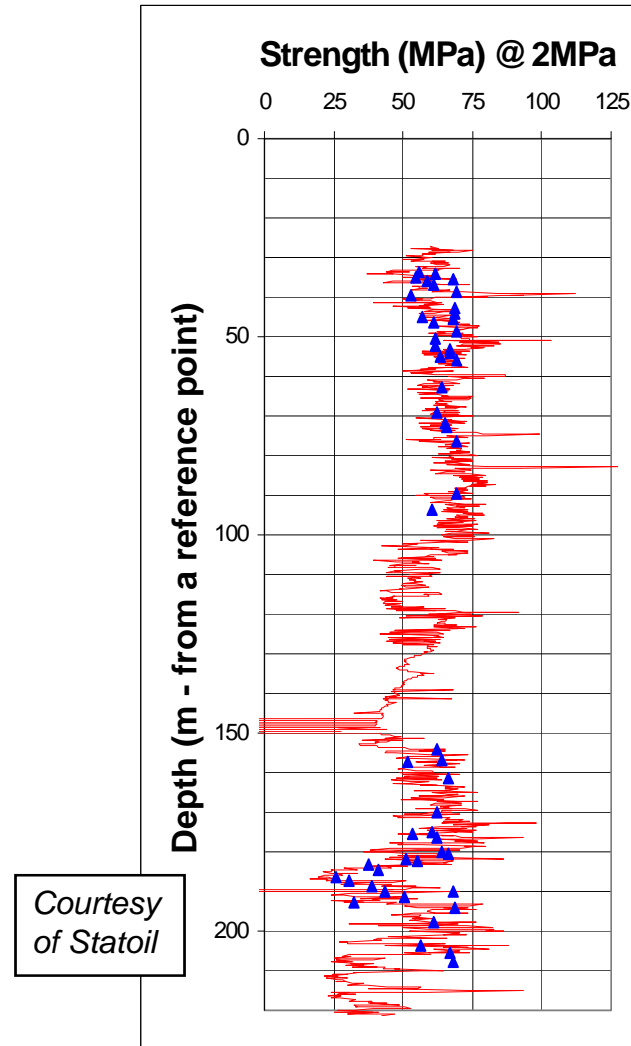
Strength



Simulates rock mechanical test on fictitious core

... an example:

- Prediction from logs
- ▲ Core measurements

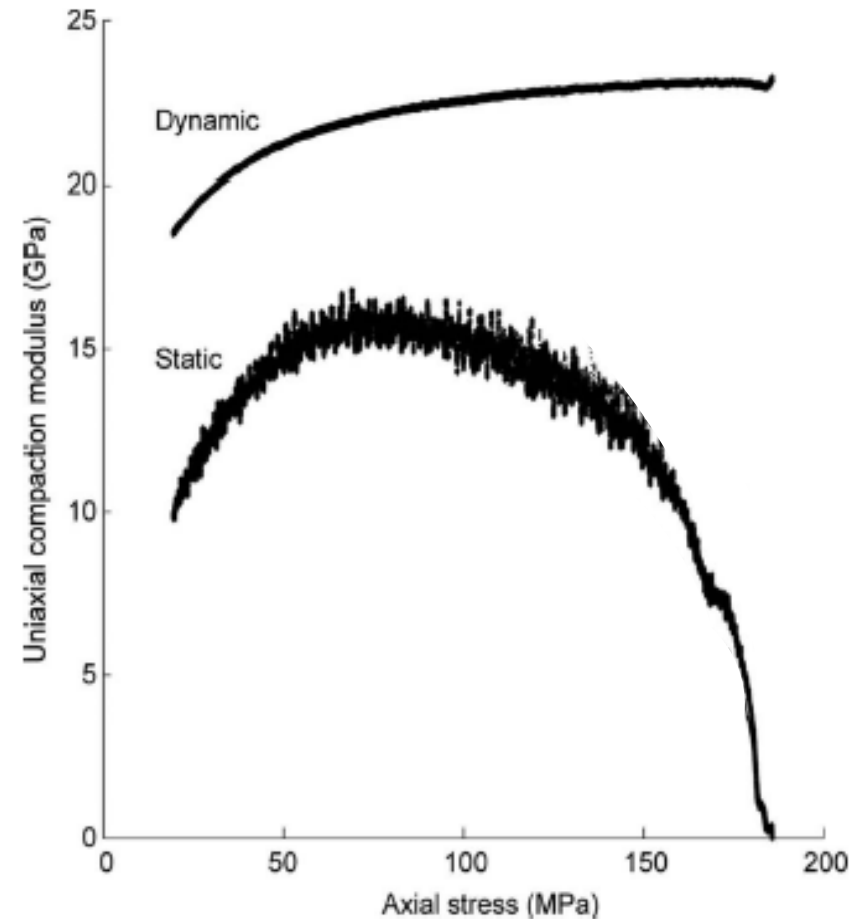


Static and dynamic moduli of soft rocks are different.

The difference changes along the stress path.

Potential causes for the difference between static and dynamic moduli:

- Strain rate
- Length of stress path
- Stress history
- Rock volume involved
- Drainage conditions
- Anisotropy



1. Stress path:

Uniaxial strain (K_0)

Static modulus = slope of stress-strain curve:

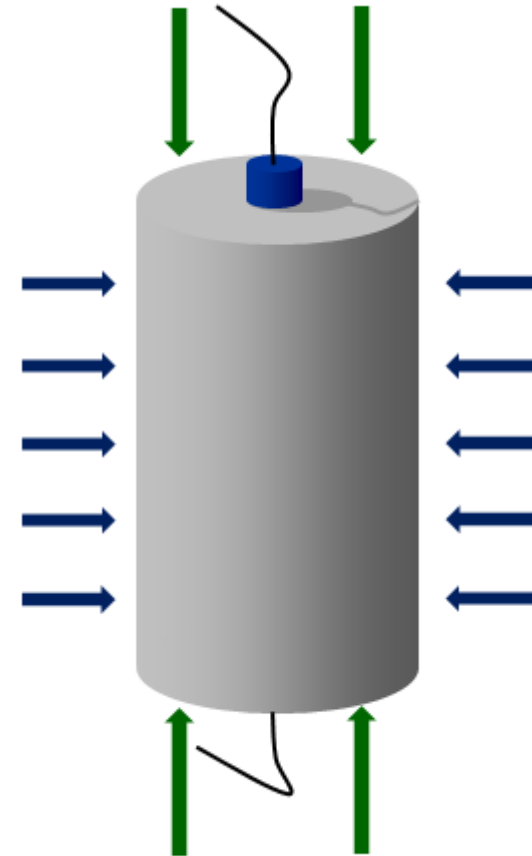
$$H = \frac{\Delta\sigma_z}{\Delta\varepsilon_z}$$

Dynamic modulus given by axial P-wave velocity:

$$H_e = \rho V_P^2$$

Potential causes for the difference between static and dynamic moduli:

- Strain rate
- Length of stress path
- Stress history
- ~~Rock volume involved~~
- Drainage conditions
- ~~Anisotropy~~



2. Saturation:

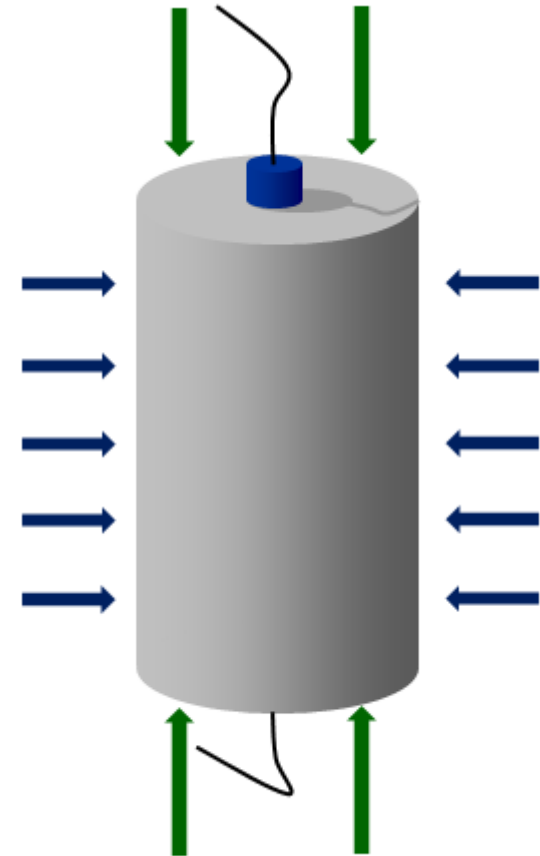
Static modulus: usually drained

Dynamic modulus: undrained

Irrelevant for dry rocks

Potential causes for the difference between static and dynamic moduli:

- Strain rate
- Length of stress path
- Stress history
- ~~Rock volume involved~~
- ~~Drainage conditions~~
- ~~Anisotropy~~

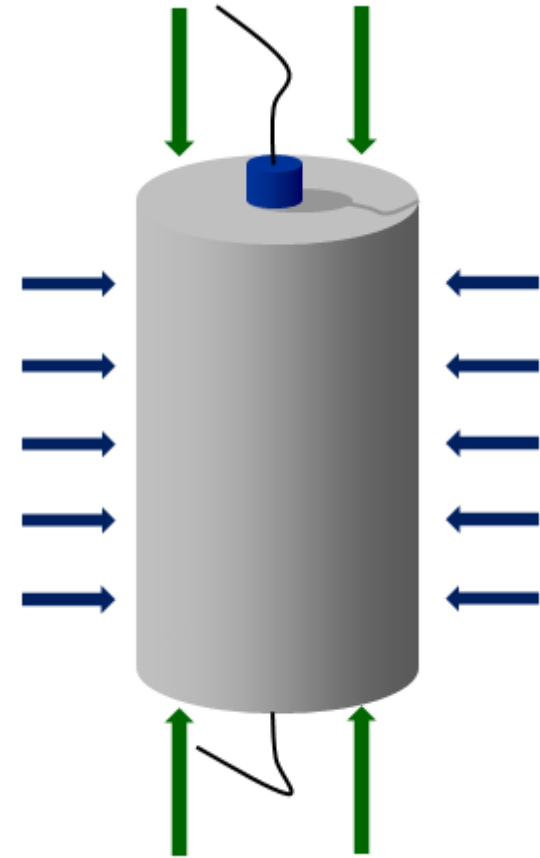


3. Loading direction:

Friction controlled shear sliding of closed cracks
- the only non-elastic process active during unloading

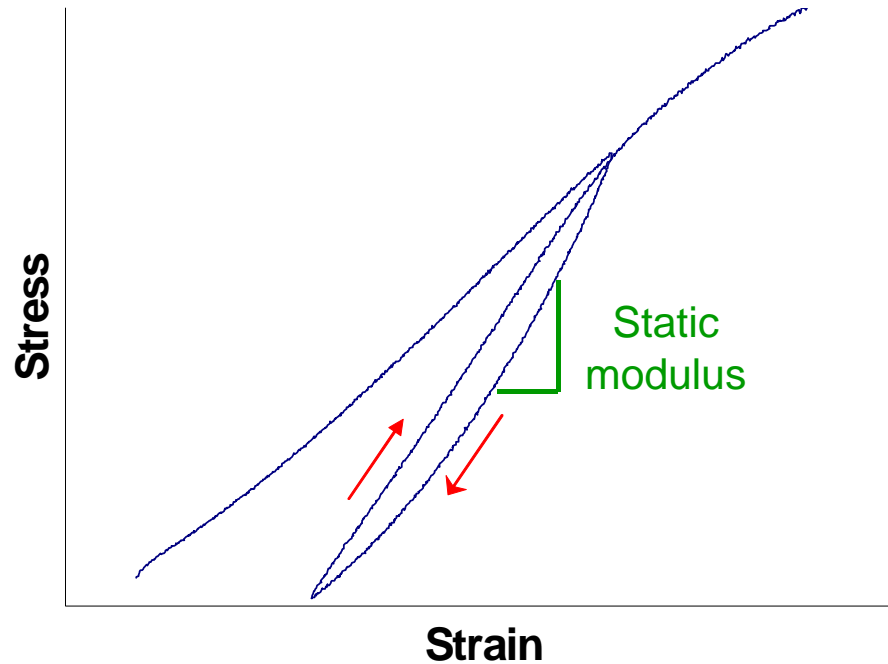
Potential causes for the difference between static and dynamic moduli:

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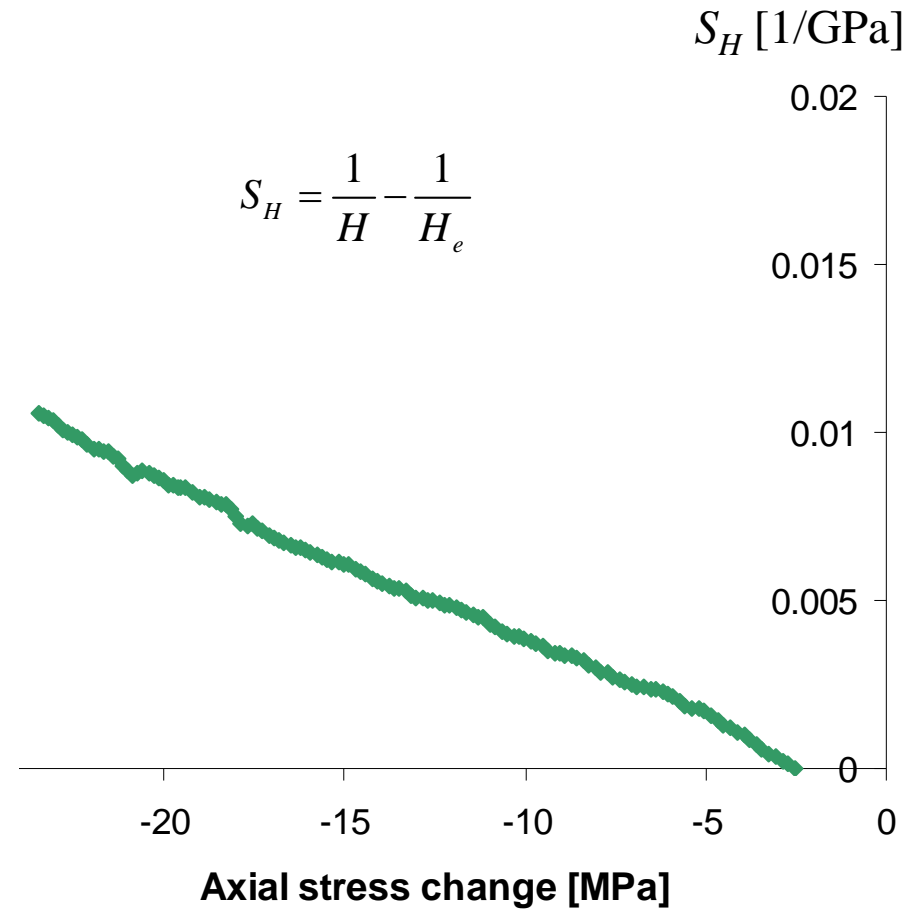


Fjær et al. (2011):

Consider the slope of the stress-strain curve during unloading



$$S_H = \frac{1}{H} - \frac{1}{H_e}$$



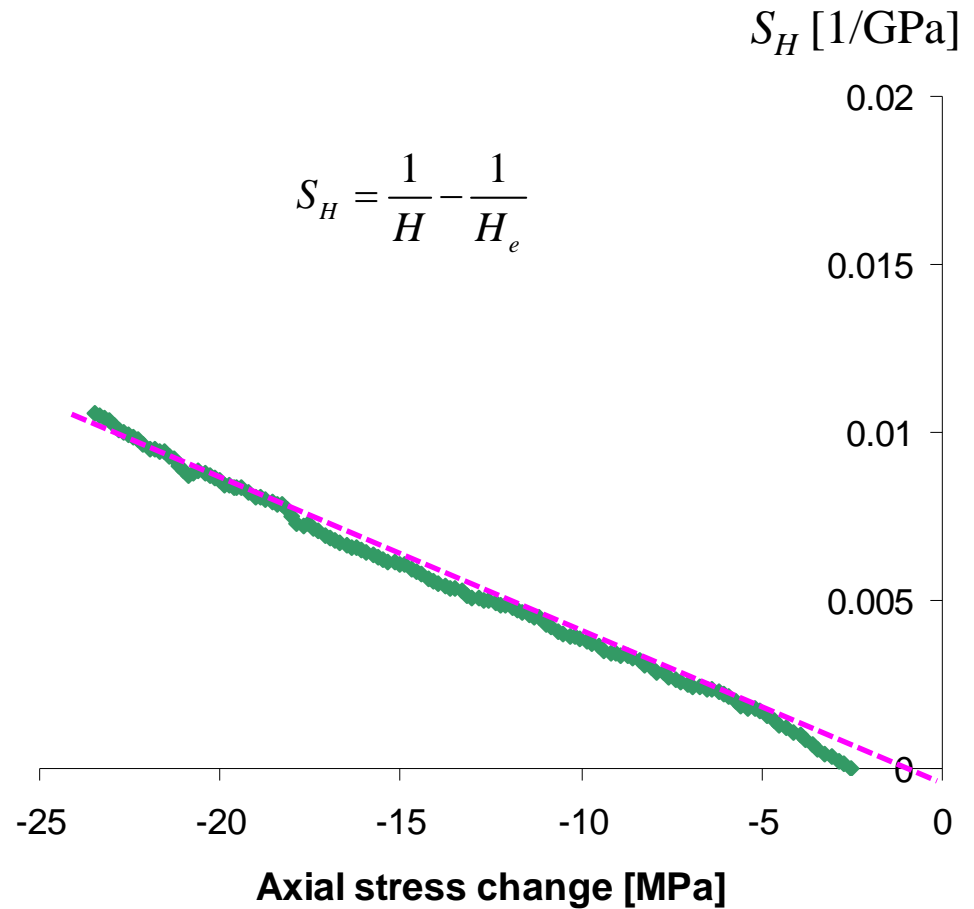
Observation:

Non-elastic compliance $S_H = \frac{1}{H} - \frac{1}{H_e}$
 increases linearly with decreasing stress

Assumption:

Linear extrapolation towards the beginning of the unloading path provides estimate of behavior for vanishing length of stress path

$$S_H = \frac{1}{H} - \frac{1}{H_e}$$



Observation:

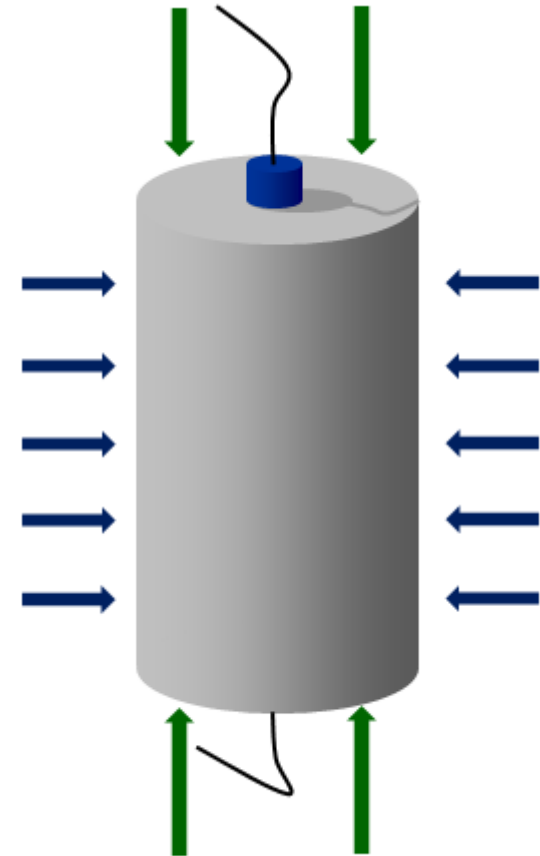
Non-elastic compliance $S_H = \frac{1}{H} - \frac{1}{H_e}$
increases linearly with decreasing stress

4. Loading direction:

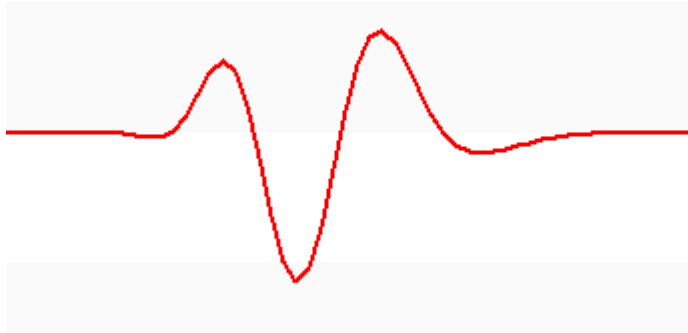
Linear extrapolation towards the beginning of the unloading path provides estimate of behavior for vanishing length of stress path

Potential causes for the difference between static and dynamic moduli:

- Strain rate
- ~~Length of stress path~~
- ~~Stress history~~
- ~~Rock volume involved~~
- ~~Drainage conditions~~
- ~~Anisotropy~~



Strain rate



Average strain rate for dynamic measurements:

$$\langle \dot{\varepsilon} \rangle = 4 f \Delta \varepsilon$$

$$\Delta \varepsilon \sim 10^{-7} \quad = \text{strain amplitude}$$

$$f = 5 \cdot 10^5 \text{ Hz} \quad = \text{frequency}$$

$$\Rightarrow \langle \dot{\varepsilon} \rangle \approx 10^{-1} \text{ s}^{-1}$$

For "static" deformations

$$\dot{\varepsilon} \approx 10^{-6} \text{ s}^{-1}$$

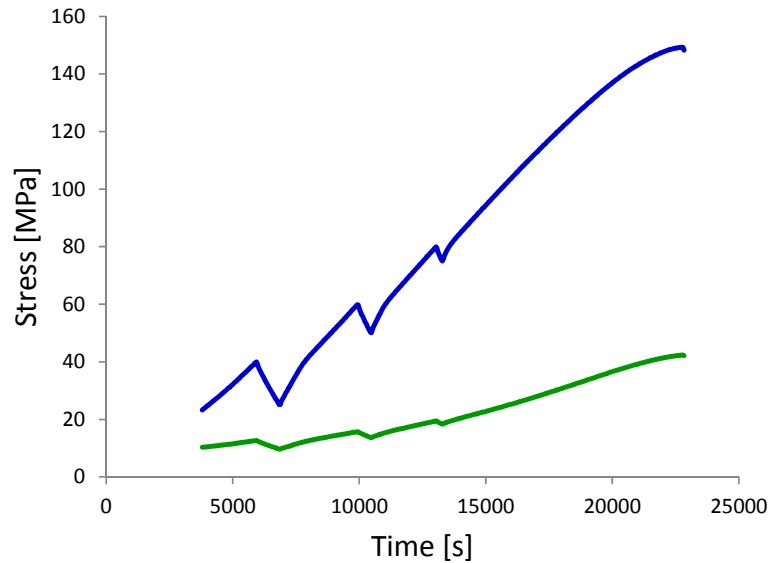
Corresponds to an acoustic wave with $f \approx 1 \text{ Hz}$

If strain rate is the only cause for the difference between the static and dynamic moduli, then

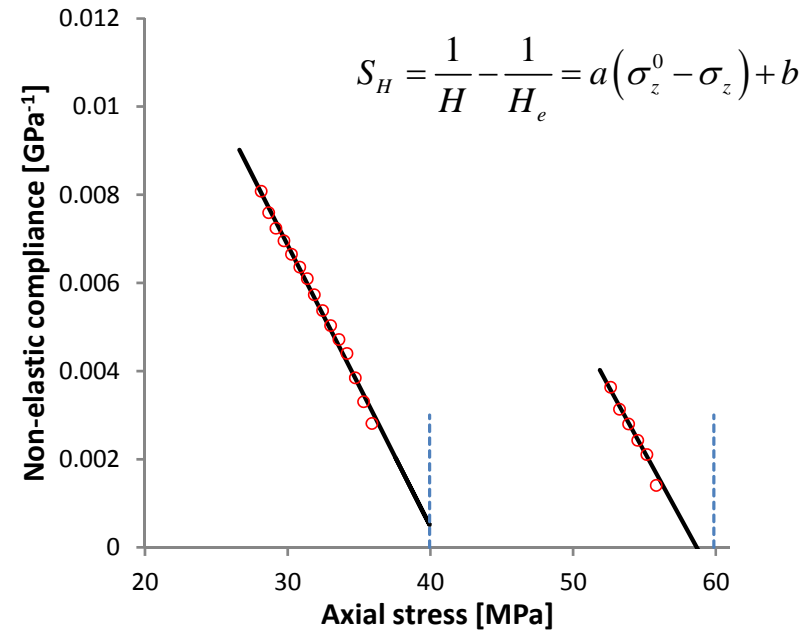
static modulus ↔ dynamic modulus at $\approx 1 \text{ Hz}$

Castlegate sandstone

Dry, clay free – presumably no significant dispersion



- K0 stress path
- Unloading sequences



σ_z MPa	σ_r MPa	b 10^{-4} GPa^{-1}	$V_{Ph} - V_{Pl}$ m/s	V_{Ph} m/s
40	13	5.2 ± 1.8	16 ± 36	3222
60	16	-7.2 ± 2.0	-24 ± 38	3285

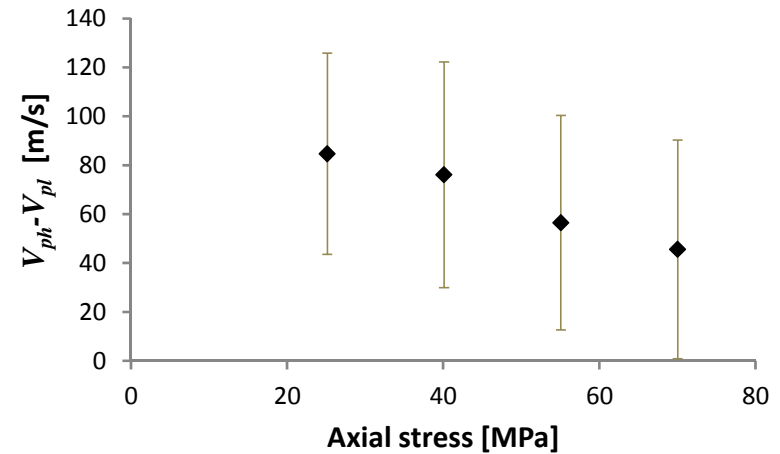
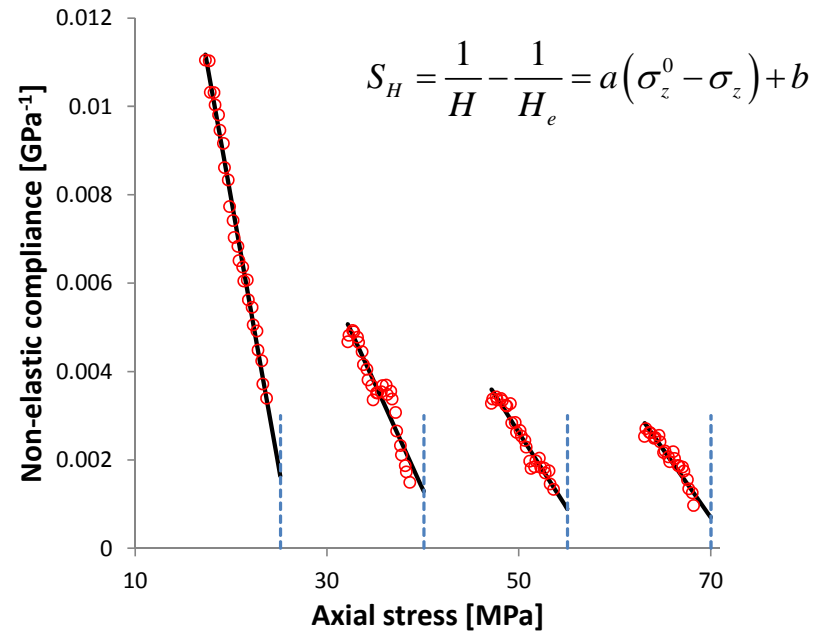
No measurable dispersion

Uncertainties: $\Delta b, \Delta \rho = 1\%$, $\Delta V_{Ph} = 1\%$

Fjær et al. (2012):

Berea sandstone

Dry, 8% clay – possibly some dispersion



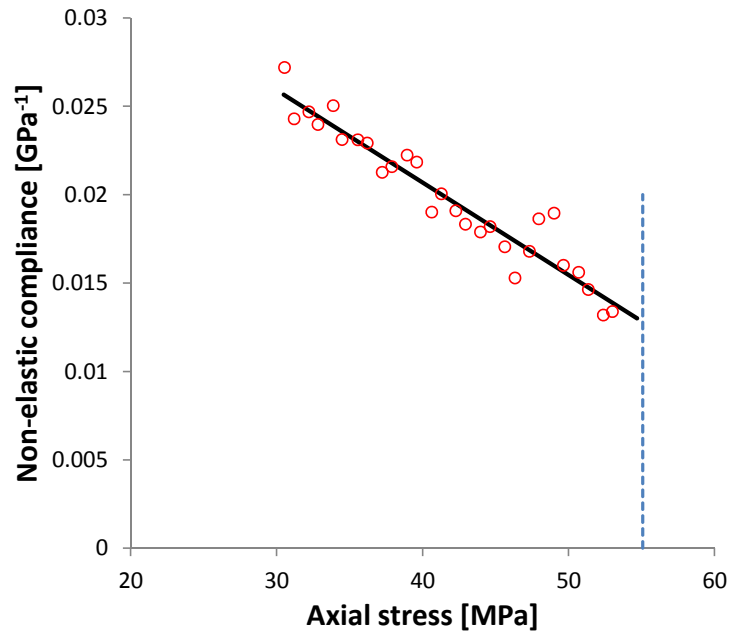
Uncertainties: Δb , $\Delta \rho = 1\%$, $\Delta V_{ph} = 1\%$

σ_z MPa	σ_r MPa	b 10^{-4} GPa^{-1}	$V_{Ph} - V_{Pl}$ m/s	V_{Ph} m/s
25	11	16.1 ± 2.8	85 ± 41	3703
40	13	12.8 ± 3.6	76 ± 46	3853
55	15	8.8 ± 1.9	56 ± 44	3929
70	17	6.9 ± 2.1	46 ± 45	3963

Significant, measurable dispersion,
decreasing with increasing stress

Mancos shale

Saturated, with 13% illite/smectite, 5% kaolinite, etc.
 – probably significant dispersion



σ_z MPa	σ_r MPa	b 10^{-4} GPa^{-1}	$V_{Ph} - V_{Pl}$ m/s	V_{Ph} m/s
55	18	129 ± 12	846 ± 112	4163

Uncertainties: $\Delta b, \Delta \rho = 1\%$, $\Delta V_{Ph} = 1\%$

Significant dispersion,
 far beyond the resolution limit
 for the method

Other applications

We measure,

- the axial P-wave velocity $V_{Pz} = \sqrt{\frac{C_{33}^e}{\rho}}$

- the axial S-wave velocity $V_{S zr} = \sqrt{\frac{C_{44}^e}{\rho}}$

- the axial stress σ_z

- the radial stress σ_r

We want

Thomsen's δ $\delta = \frac{(C_{13}^e + C_{44}^e)^2 - (C_{33}^e - C_{44}^e)^2}{2C_{33}^e (C_{33}^e - C_{44}^e)}$

$$r_{44}^e \equiv \frac{C_{44}^e}{C_{33}^e} = \left(\frac{V_{S, zr}}{V_{P, z}} \right)^2$$

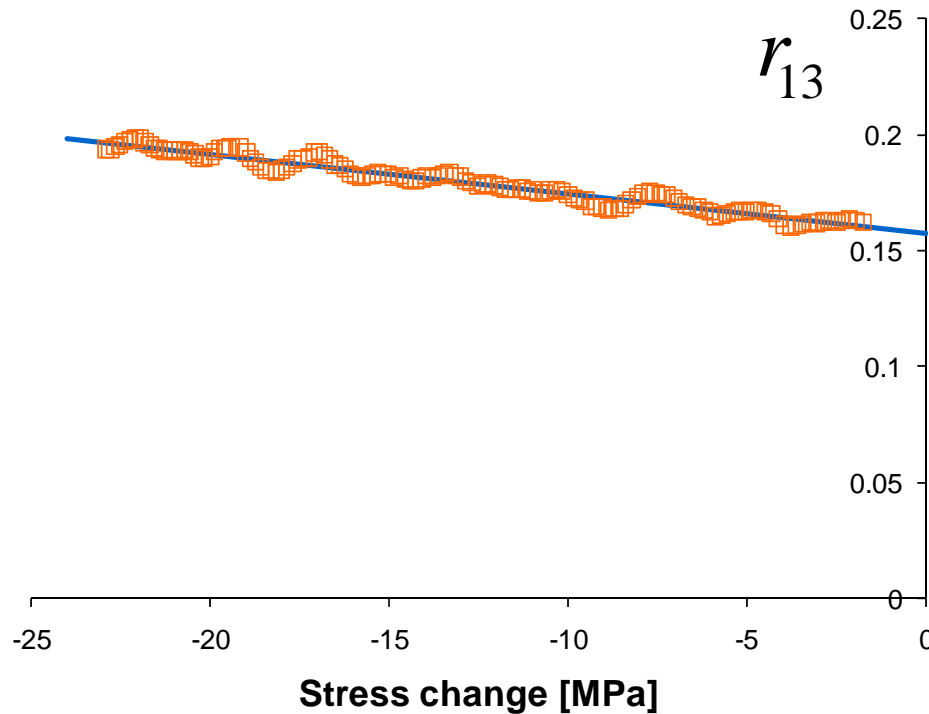
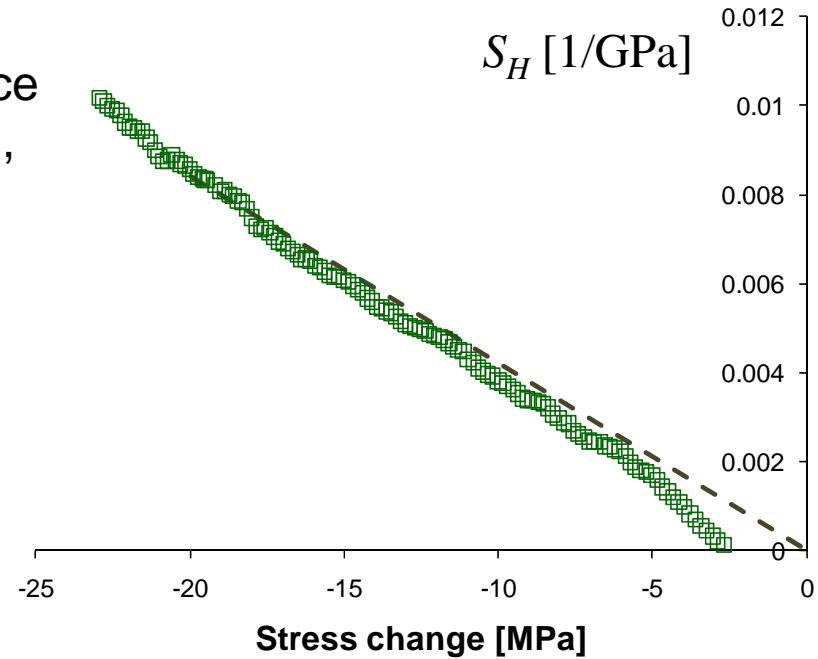
$$r_{13} \equiv \frac{C_{13}}{C_{33}} = \frac{\Delta\sigma_r}{\Delta\sigma_z}$$

$$\delta = \frac{(r_{13}^e + r_{44}^e)^2 - (1 - r_{44}^e)^2}{2(1 - r_{44}^e)}$$

\Rightarrow We have r_{44}^e & r_{13} , we need r_{13}^e

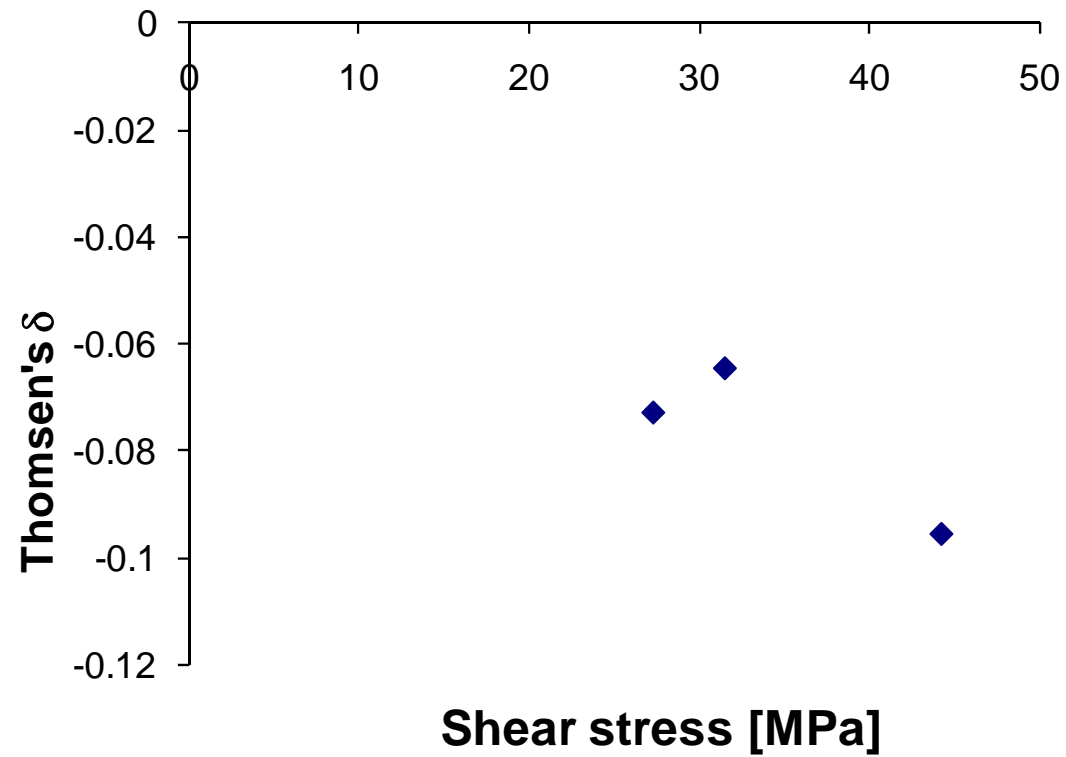
We have seen that the non-elastic compliance vanish at the turning point on the stress path, (the approach is linear during unloading)

- which means:
static modulus \rightarrow dynamic modulus
at the turning point



Therefore –
We assume that
static $r_{13} \rightarrow$ dynamic r_{13}
at the turning point

Example (dry Castlegate sandstone):



References:

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Fjær, E. (1999) "Static and dynamic moduli of weak sandstones". In *Proceedings of the 37th U.S. Rock Mechanics Symposium*, eds Amadei et al., 675-681

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