

Crack models

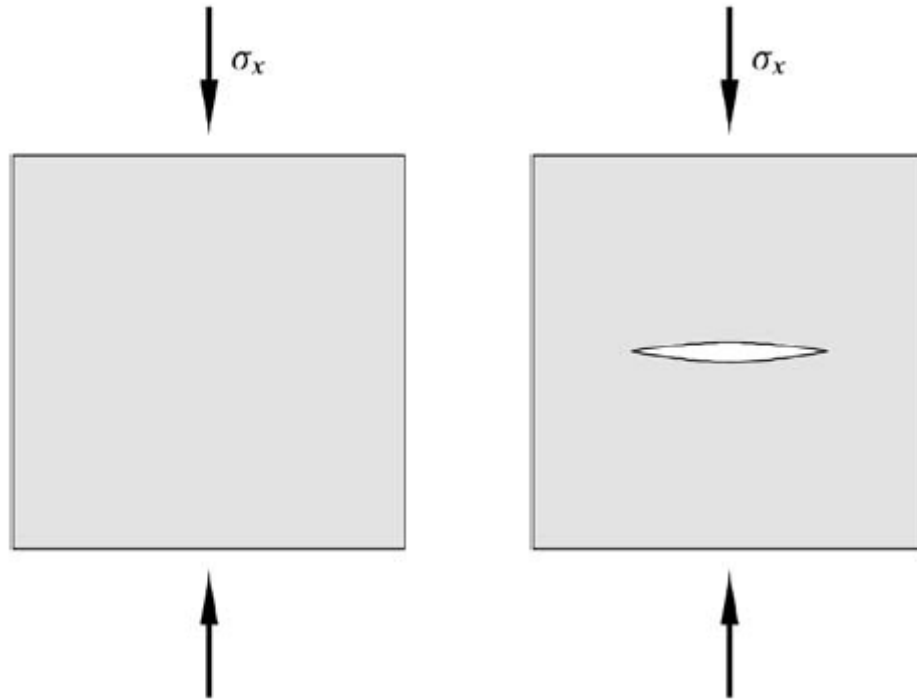
ROSE

Rock Physics and Geomechanics

Course 2012

Erling Tjørrer

Cracks have a strong impact on rock behavior

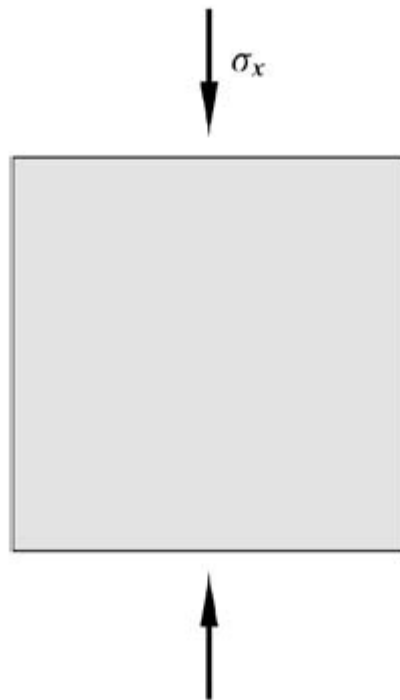


The presence of cracks
reduces the stiffness
of the rock

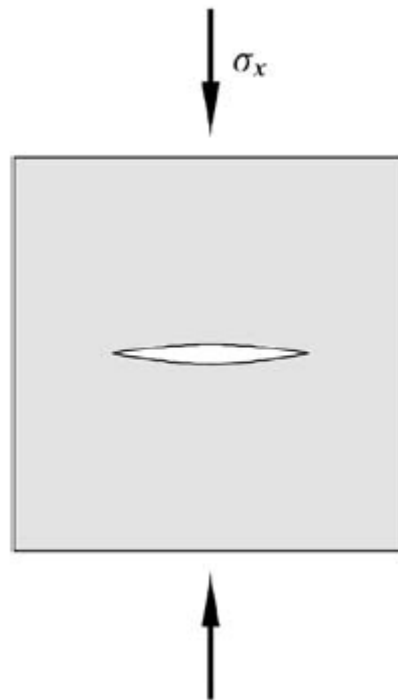
$$E = \frac{\Delta\sigma_x}{\Delta\varepsilon_x}$$

$$E_{\text{eff}} = \frac{\Delta\sigma_x}{\Delta\varepsilon_x} = E(1 - Q\xi)$$

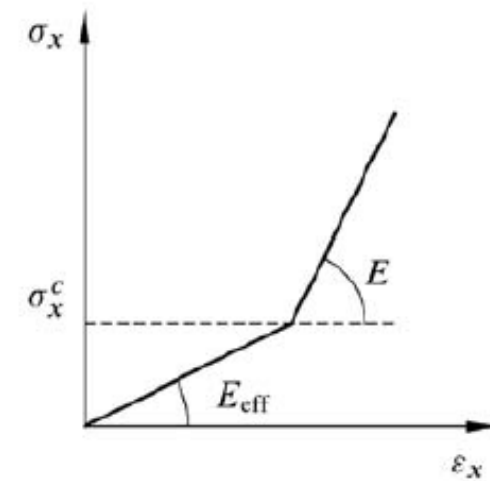
Cracks have a strong impact on rock behavior



$$E = \frac{\Delta\sigma_x}{\Delta\varepsilon_x}$$

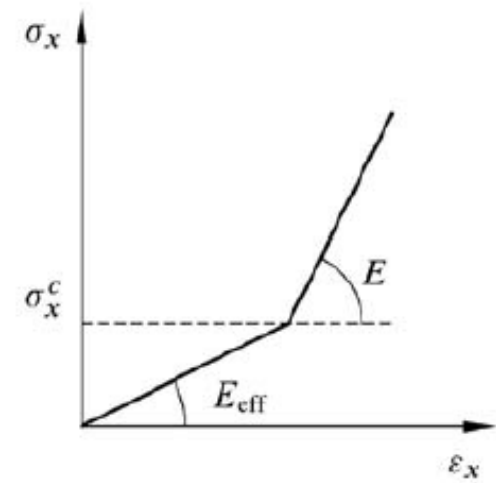
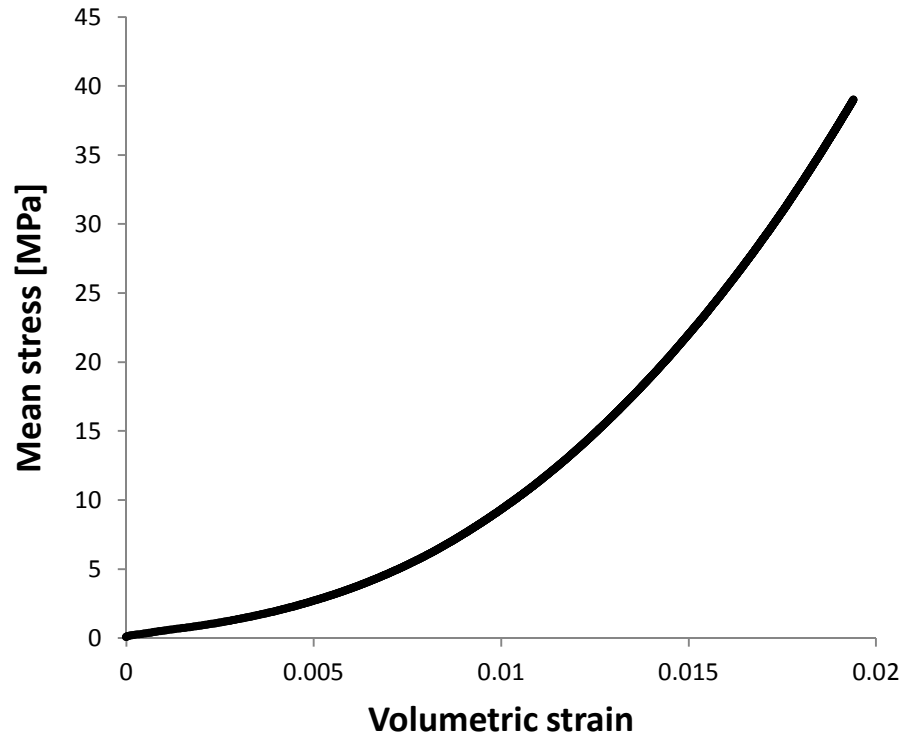


$$E_{\text{eff}} = \frac{\Delta\sigma_x}{\Delta\varepsilon_x} = E(1 - Q\xi)$$

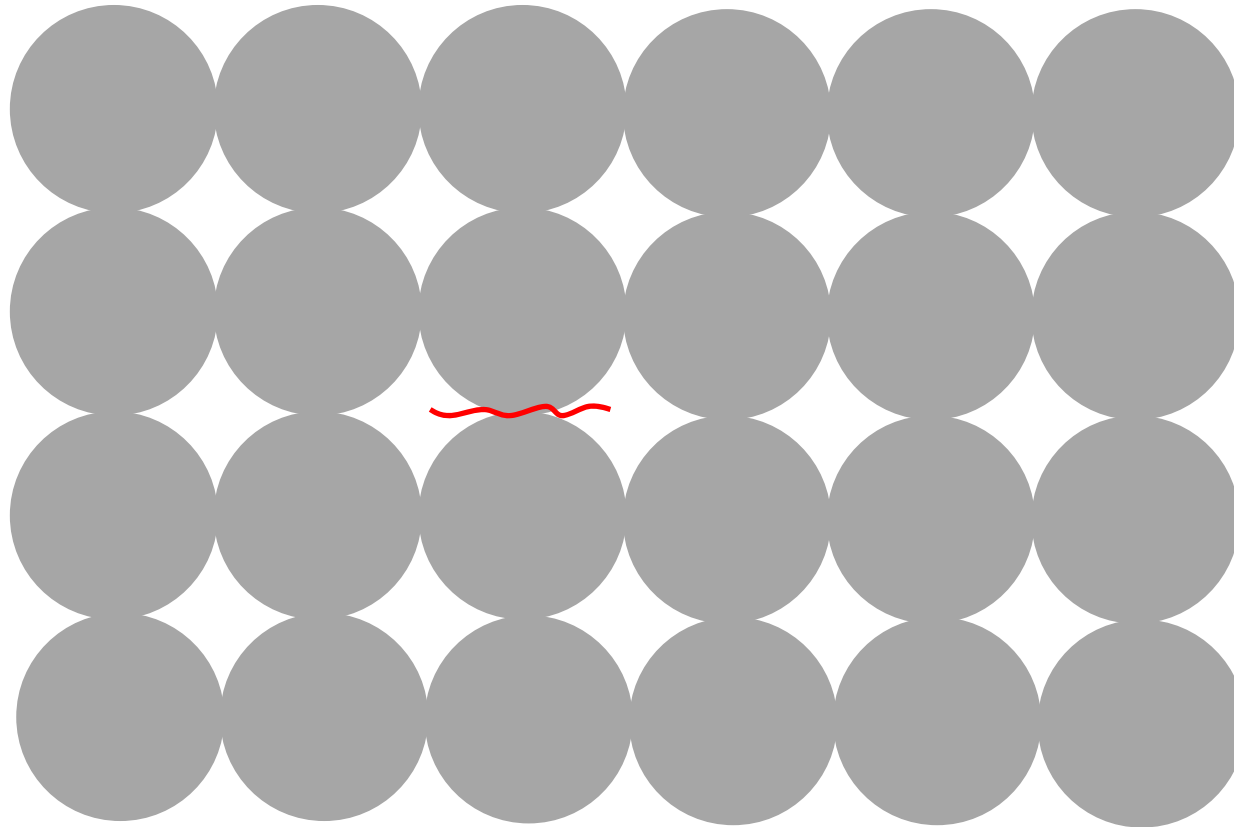


Increasing stress
 ⇒ Crack closure
 ⇒ Increasing stiffness

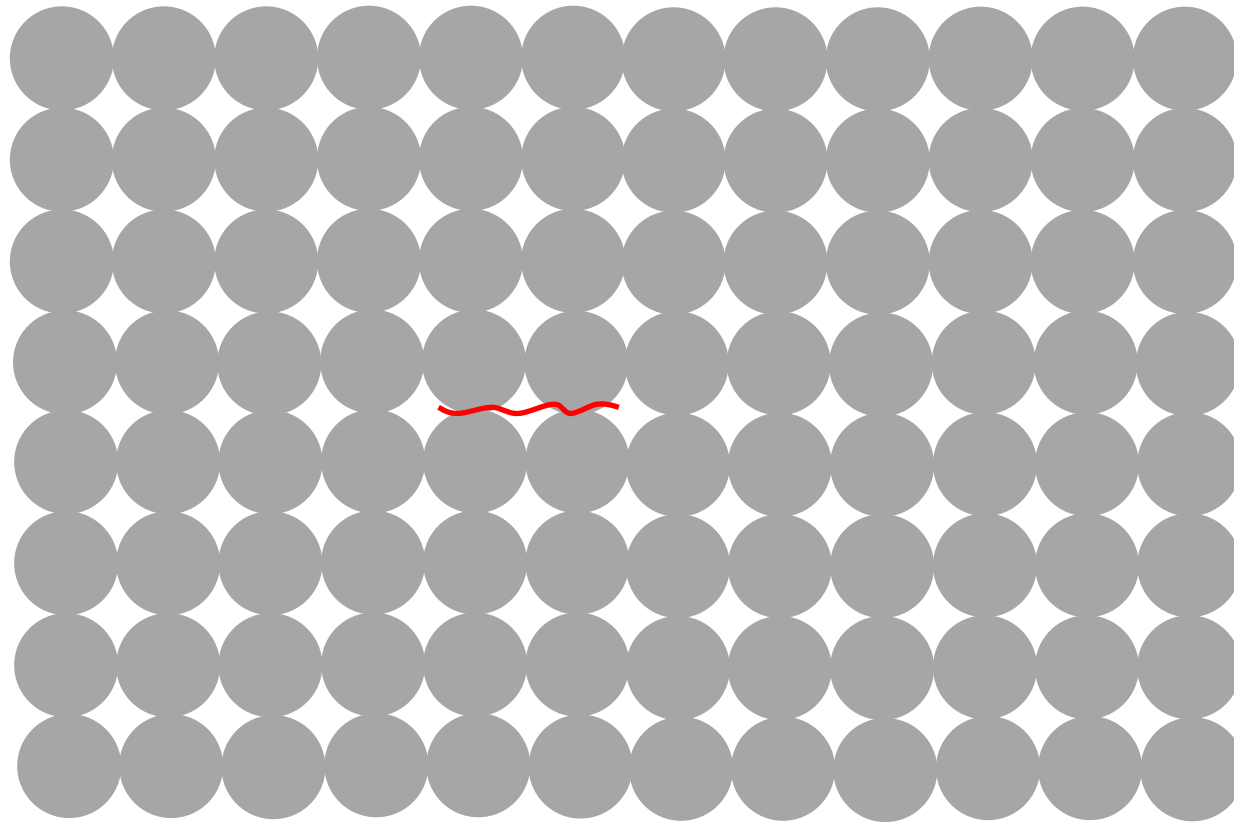
Cracks have a strong impact on rock behavior



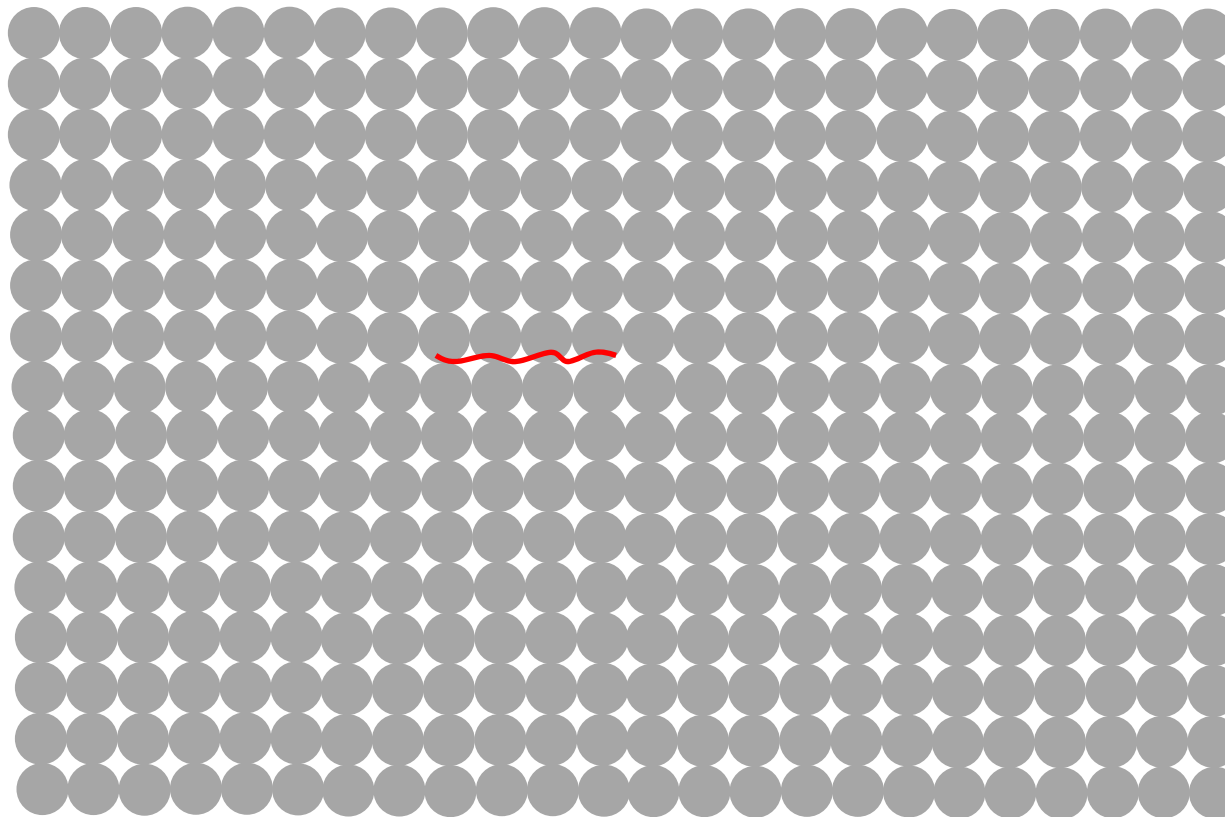
Increasing stress
⇒ Crack closure
⇒ Increasing stiffness



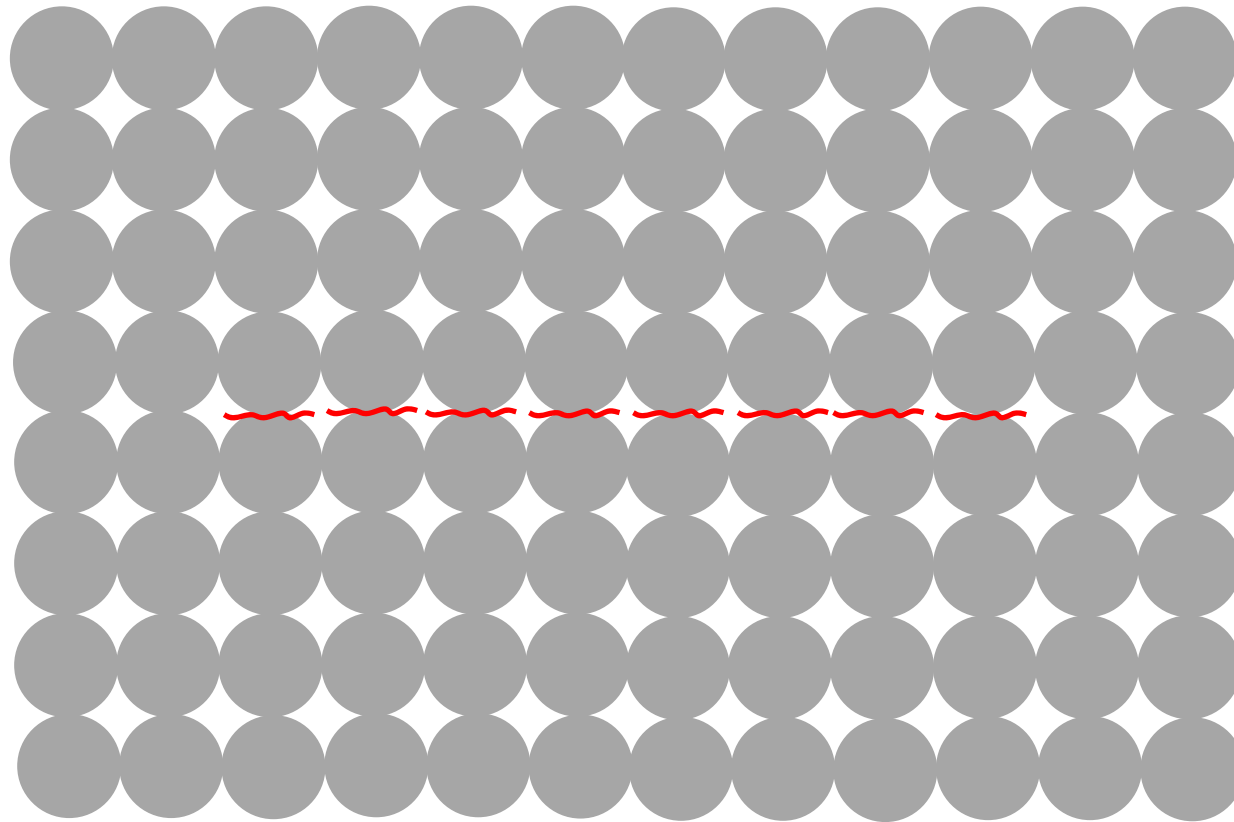
Crack – or failed grain contact?



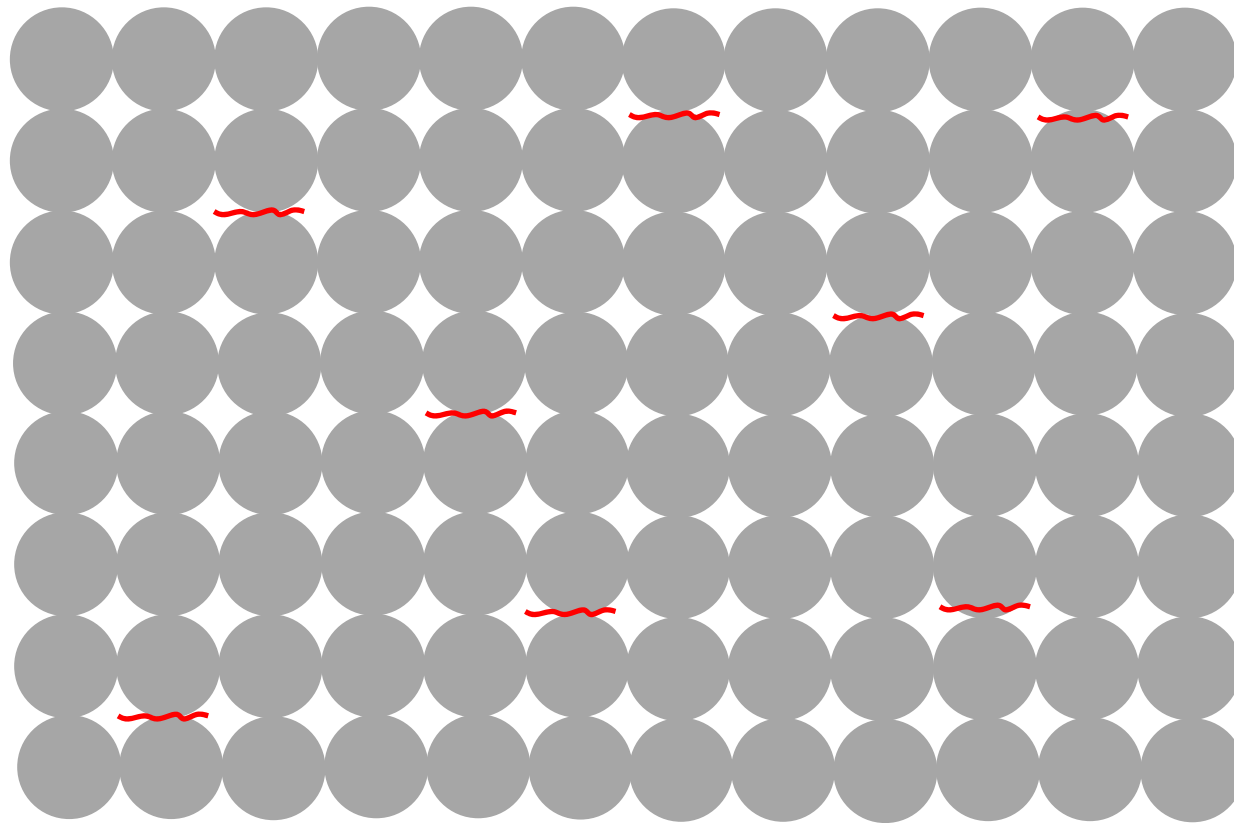
Crack – or failed grain contacts?



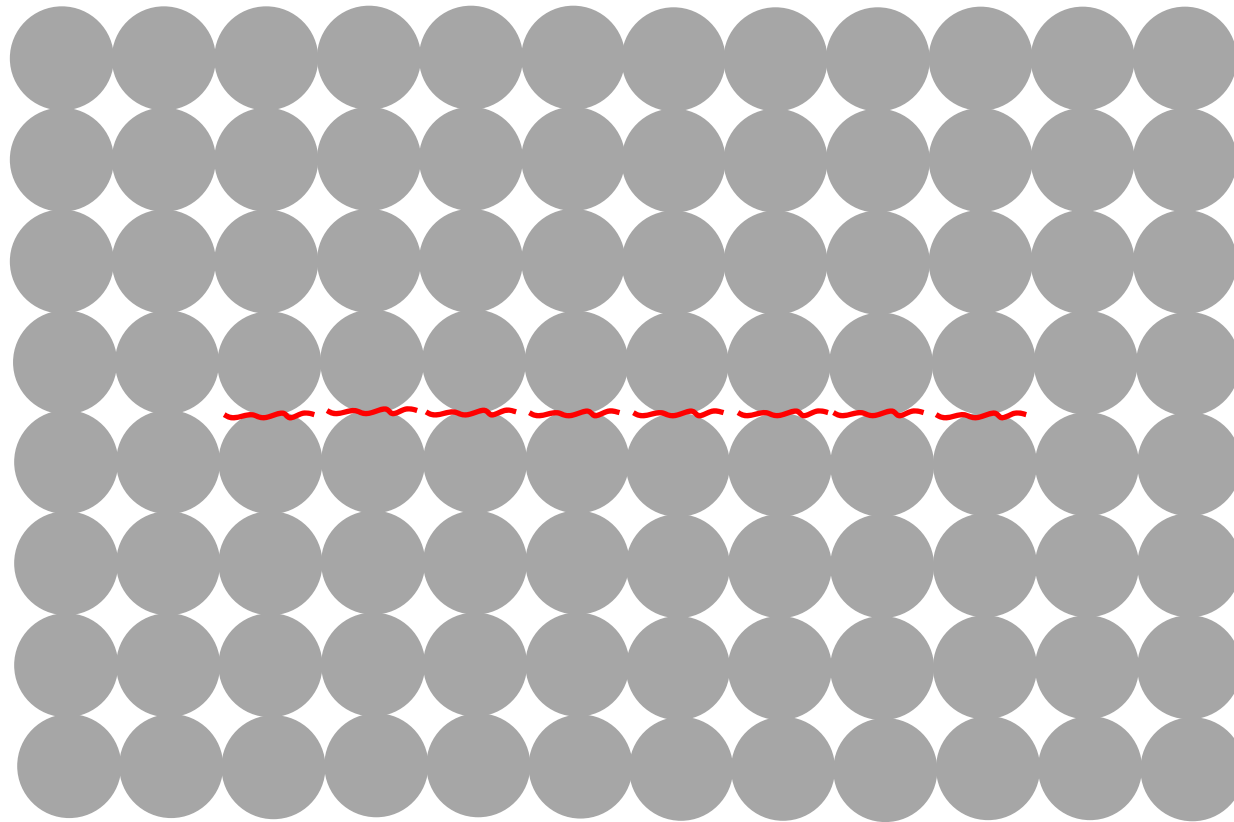
Crack – or failed grain contacts?



One large crack –



– or several failed grain contacts?



An assembly of failed grain contacts = a large crack
has a much stronger impact on rock stiffness
than the sum of the individual failed grain contacts

Open, flat cracks

Crack density: $\xi = n \langle a^3 \rangle$

n = number of cracks per unit volume



Isotropic distribution of cracks:

$$K^* = K_s \left(1 - \frac{16}{9} \frac{1 - \nu_s^2}{1 - 2\nu_s} D \xi \right)$$

$$G^* = G_s \left(1 - \frac{32}{45} (1 - \nu_s) \left[D + \frac{3}{2 - \nu_s} \right] \xi \right)$$

Not consistent with Biot, but
 – we may use the model to predict
 the properties of the dry material,
 which gives us the *frame moduli*.

Drainage
parameter

$$\frac{1}{D} = 1 + \frac{4}{3\pi\gamma} \frac{1 - \nu_s^2}{1 - 2\nu_s} \frac{K_f}{K_s}$$

Flat cracks: $\gamma = c/a \rightarrow 0$

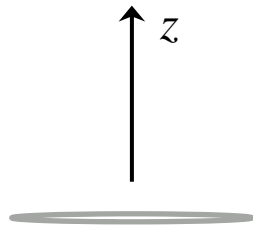
\Rightarrow

$D \rightarrow 0$ for saturated rock

$D = 1$ for dry rock

Non-isotropic distribution of cracks \Rightarrow **anisotropy**

$$C_{ij}^* = C_{ij}^o \left(1 - \sum_k Q_{ij}^k \zeta_k \right)$$



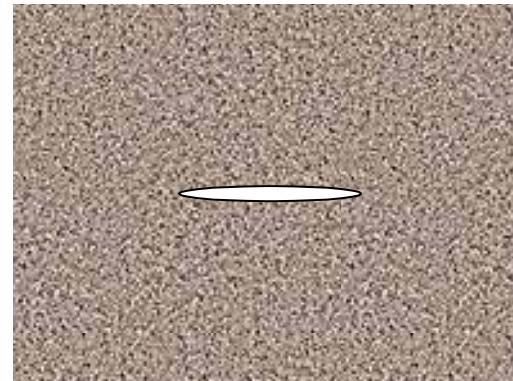
$$Q^k = \frac{16}{3} \begin{bmatrix} \frac{\nu_s^2}{1-2\nu_s} D & \frac{\nu_s(1-\nu_s)}{1-2\nu_s} D & \frac{(1-\nu_s)^2}{1-2\nu_s} D & 0 & 0 & 0 \\ \frac{\nu_s(1-\nu_s)}{1-2\nu_s} D & \frac{\nu_s^2}{1-2\nu_s} D & \frac{(1-\nu_s)^2}{1-2\nu_s} D & 0 & 0 & 0 \\ \frac{(1-\nu_s)^2}{1-2\nu_s} D & \frac{(1-\nu_s)^2}{1-2\nu_s} D & \frac{(1-\nu_s)^2}{1-2\nu_s} D & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{1-\nu_s}{2-\nu_s} & 0 & 0 \\ 0 & 0 & 0 & 0 & \frac{1-\nu_s}{2-\nu_s} & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix}$$

Non-isotropic distribution of cracks \Rightarrow **anisotropy**

Cracks control the velocities:





















Large reduction
in velocity
and amplitude



No effect (almost) on velocity
and amplitude of P-wave

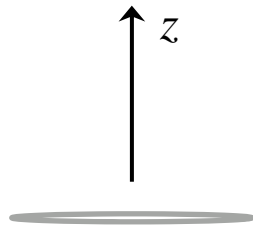
Some effect on vertically
polarized S-wave

Cracks can explain.....

	<i>Propagation</i>	<i>Polarization</i>	<i>Crack orientation</i>	<i>Velocity reduction</i>
P-wave				Very strong
P-wave				Weak
P-wave				Weak
S-wave				Strong
S-wave				Strong
S-wave				None

No drainage $\Rightarrow D = 0 \Rightarrow$ (nearly) no P-wave anisotropy

$$C_{ij}^* = C_{ij}^o \left(1 - \sum_k Q_{ij}^k \zeta_k \right)$$



$$Q^k = \frac{16}{3} \begin{bmatrix} \frac{\nu_s^2}{1-2\nu_s} D & \frac{\nu_s(1-\nu_s)}{1-2\nu_s} D & \frac{(1-\nu_s)^2}{1-2\nu_s} D & 0 & 0 & 0 \\ \frac{\nu_s(1-\nu_s)}{1-2\nu_s} D & \frac{\nu_s^2}{1-2\nu_s} D & \frac{(1-\nu_s)^2}{1-2\nu_s} D & 0 & 0 & 0 \\ \frac{(1-\nu_s)^2}{1-2\nu_s} D & \frac{(1-\nu_s)^2}{1-2\nu_s} D & \frac{(1-\nu_s)^2}{1-2\nu_s} D & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{1-\nu_s}{2-\nu_s} & 0 & 0 \\ 0 & 0 & 0 & 0 & \frac{1-\nu_s}{2-\nu_s} & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix}$$

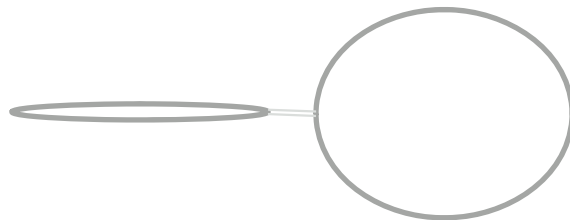
"Saturation eliminates P-wave anisotropy"

Leon Thomsen (1995):

Pore pressure equalization between cracks and pores

$$\frac{1}{D_{cp}} = 1 + \left[\frac{3}{2} \frac{1 - \nu_s}{1 - 2\nu_s} \frac{\phi_p}{\phi} + \frac{4}{3\pi\gamma} \frac{1 - \nu_s^2}{1 - 2\nu_s} \left(1 - \frac{\phi_p}{\phi} \right) \right] \frac{K_f}{K_s}$$
$$\phi - \phi_p = \frac{4}{3} \pi \gamma \xi$$

Consequence: the drainage parameter will not vanish, regardless how thin the cracks are



"Saturation *does not* eliminate P-wave anisotropy in porous and permeable rocks"

General formalism for *displacement discontinuities*
(*Sayers and Kachanov, 1995*):

$$S_{ijkl}^* = S_{ijkl}^o + \frac{1}{4}(\delta_{ik}\alpha_{jl} + \delta_{il}\alpha_{jk} + \delta_{jk}\alpha_{il} + \delta_{jl}\alpha_{ik}) + \beta_{ijkl}$$

$$\alpha_{ij} = \frac{1}{V} \sum_r B_T^r n_i^r n_j^r S_r$$

$$\beta_{ijkl} = \frac{1}{V} \sum_r (B_N^r - B_T^r) n_i^r n_j^r n_k^r n_l^r S_r$$

For open, penny-shaped cracks $B_N \approx B_T \Rightarrow \beta_{ijkl} \approx 0$

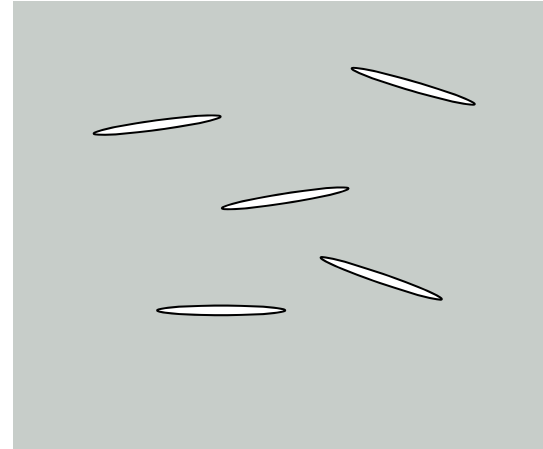
This approximation is not valid in general,
hence the open crack model is too simple.

However, we may compensate for this by allowing the
drainage parameter D be an adjustable parameter.

Many cracks \Rightarrow crack interactions

- the presence of one crack may affect the influence of another

$$K^* = K_s (1 - Q(\nu_s) \zeta)$$



Self-consistent models:

Interactions are taken into account by giving the rock around a crack the properties of the effective medium

$$K^* = K_s - K^* Q(\nu_s^*) \zeta$$

\Rightarrow

$$K^* = \frac{K_s}{1 + Q(\nu_s^*) \zeta}$$

Alternative, equally valid procedures

$$K^* = K_s (1 - Q(v_s) \zeta)$$

$$\frac{1}{K^*} = \frac{1}{K_s} + \frac{Q(v_s) \zeta}{K_s}$$

Self-consistency:

$$K^* = K_s - K^* Q(v^*) \zeta$$

$$\frac{1}{K^*} = \frac{1}{K_s} + \frac{Q(v^*) \zeta}{K^*}$$

⇓

$$K^* = \frac{K_s}{1 + Q(v^*) \zeta}$$

$$K^* = K_s (1 - Q(v^*) \zeta)$$

$$K^* > 0 \text{ always}$$

$$K^* = 0 \text{ for } \zeta = 1/Q(v^*)$$

There are many different self-consistent models,
giving different predictions – depending on the initial model.

The Differential Effective Medium (DEM) model resolves this discrepancy
by adding small numbers of cracks in many steps, and recalculating the
effective stiffness for each step (always working in the non-interacting limit)

$$dK^* = -K^* Q(v^*) d\zeta$$

This gives the same (*unique!*) solution for all initial models.

But – is it more correct because of that?

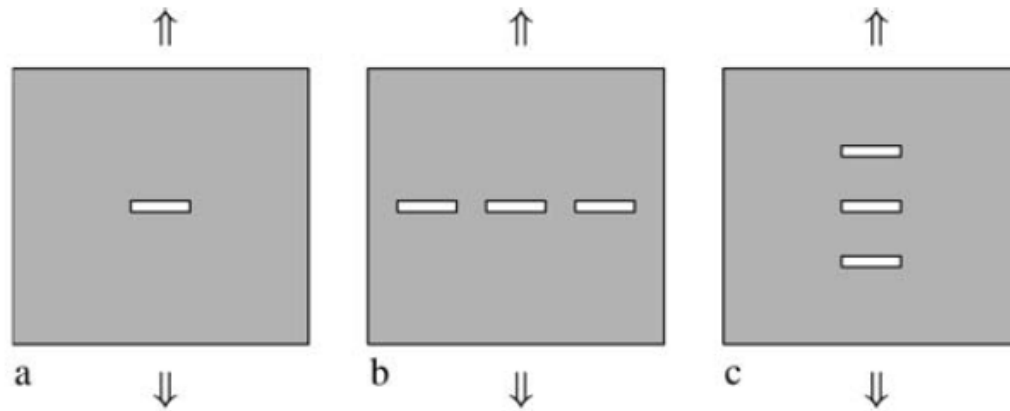
Models accounting for interactions

Many alternatives.....

- all are mathematically correct, but they give different predictions.

Which one is correct?

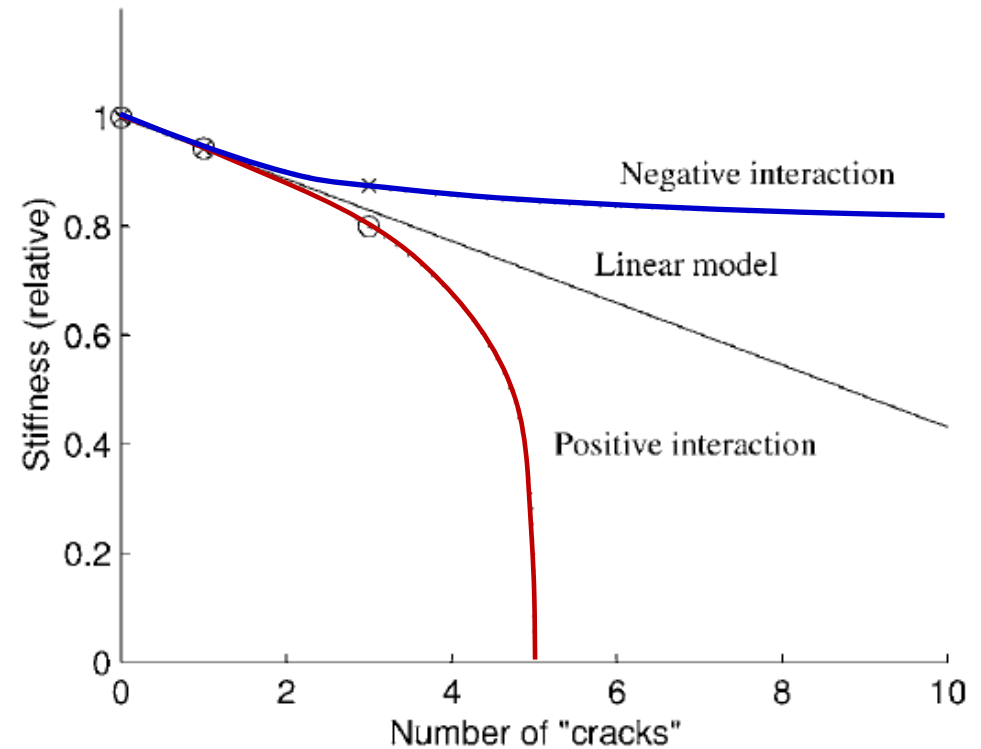
It depends....



The actual position of the cracks relative to each other determines which model is the correct one to use.

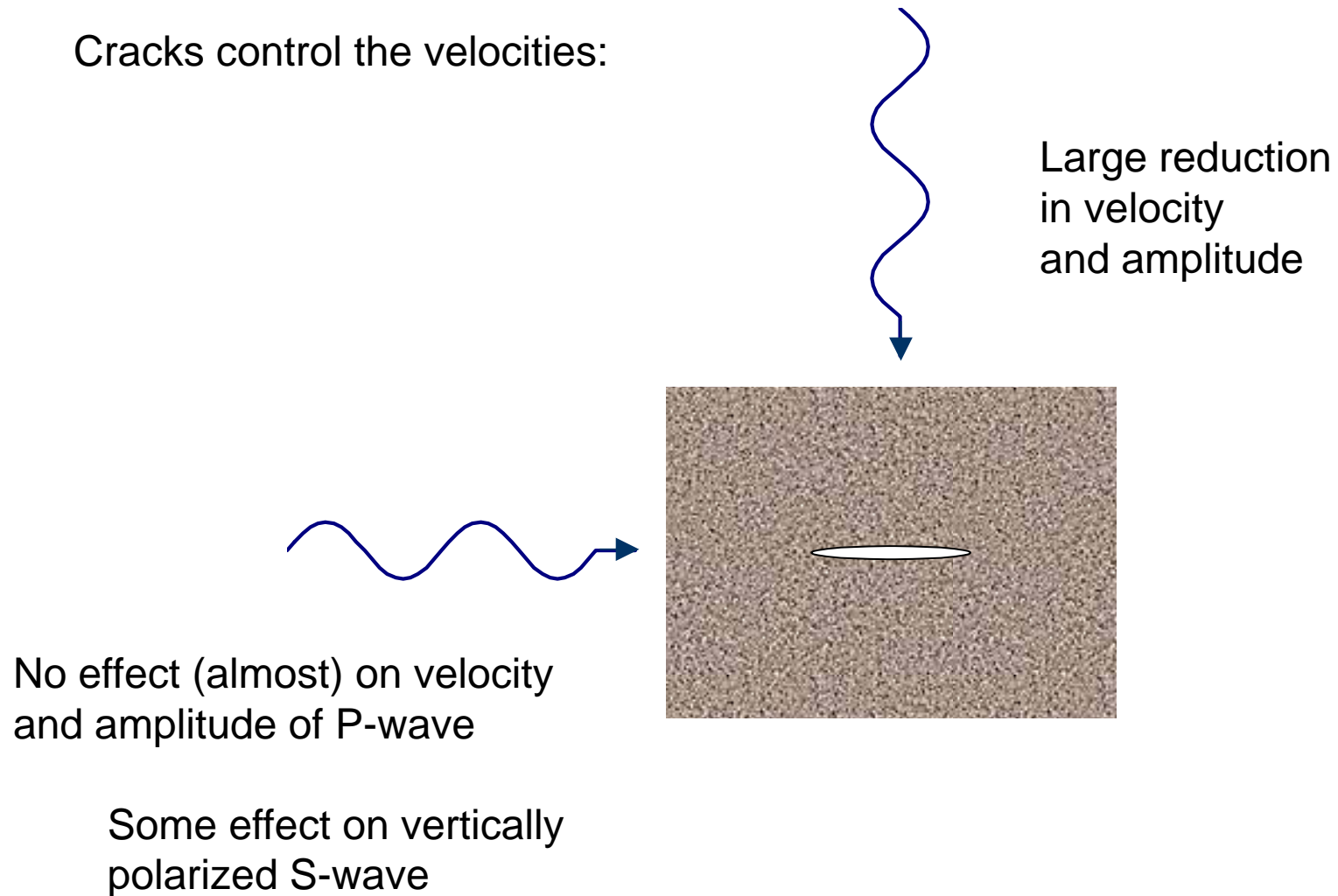
Unfortunately, no model comes with a description of how the cracks are positioned.

If we do not know how the cracks are positioned (and we usually don't), the linear, non-interacting model may be a good starting point.



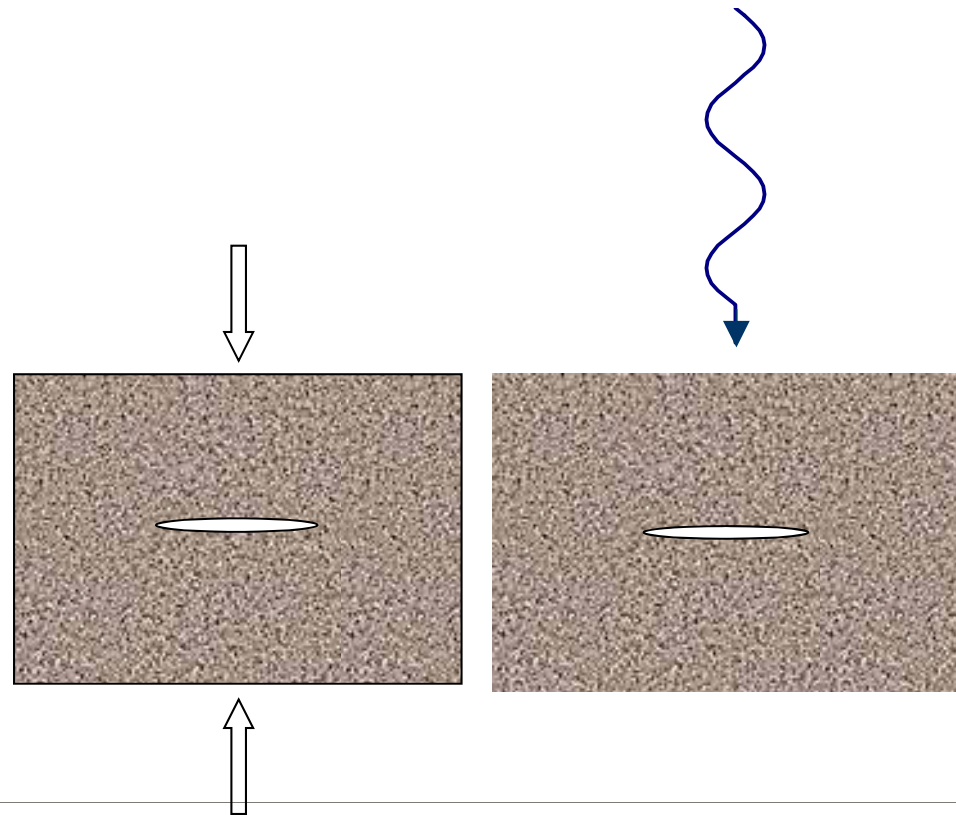
Stress effects on the rock framework

Cracks control the velocities:



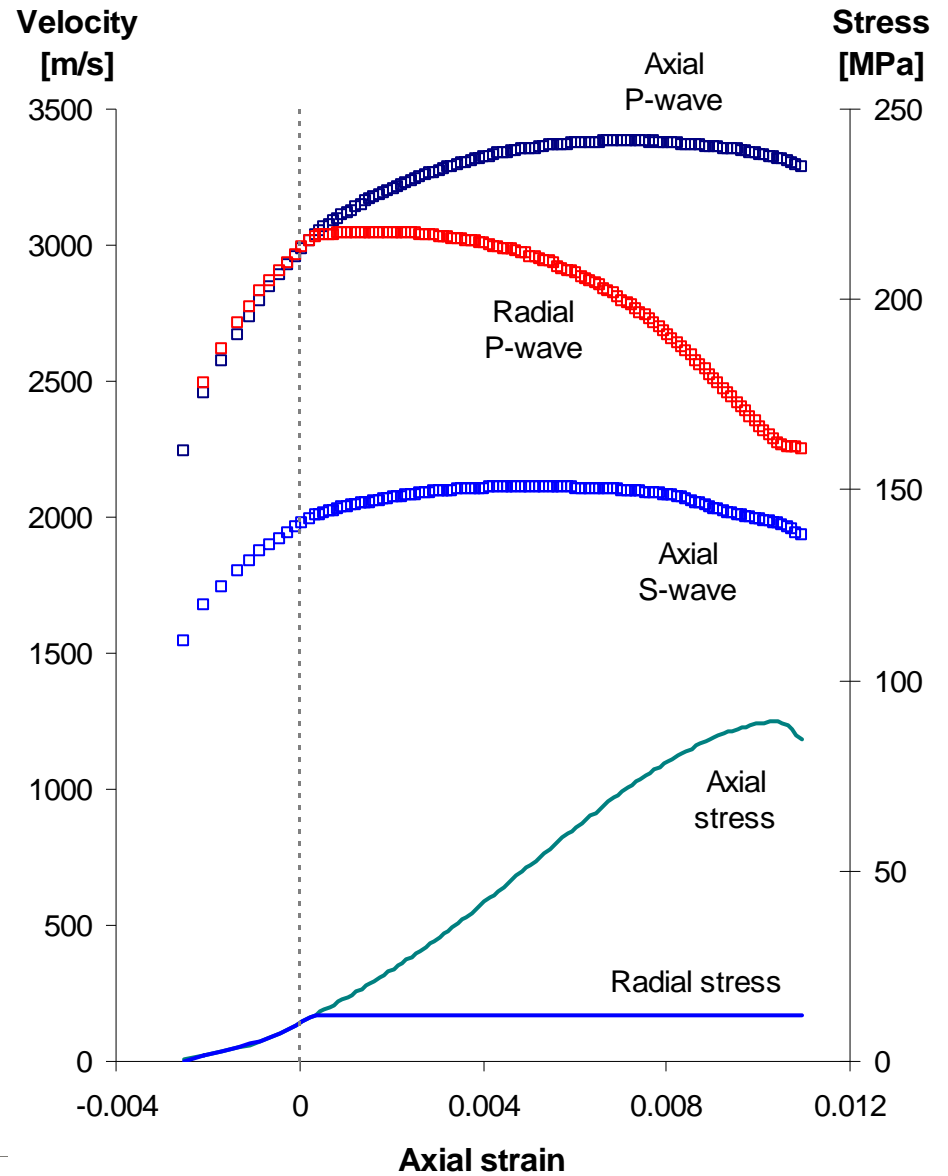
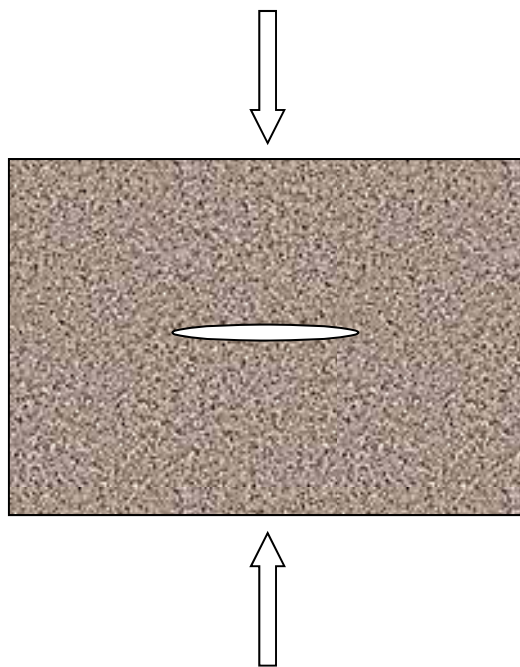
Rule of thumb:

Velocities are mostly affected by changes in the normal stress in the direction of propagation (and polarization)



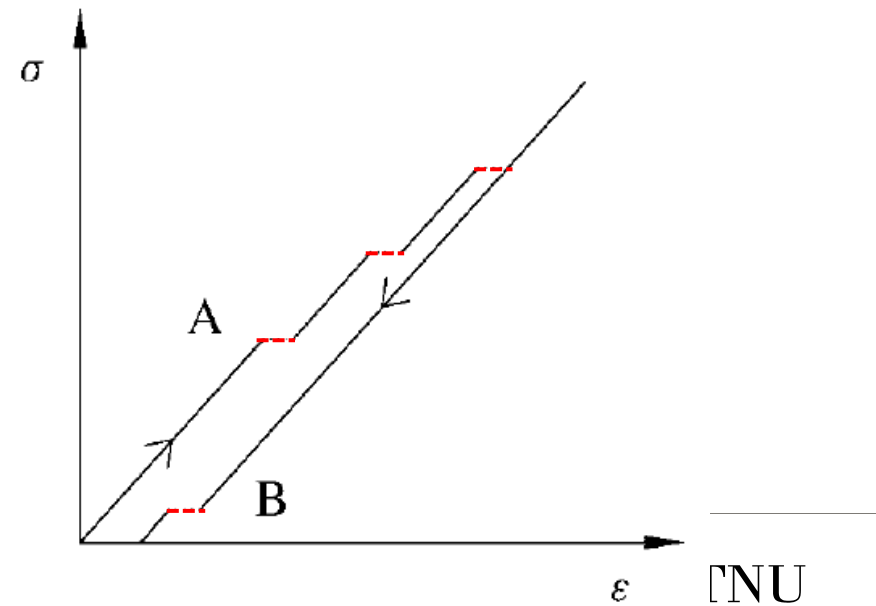
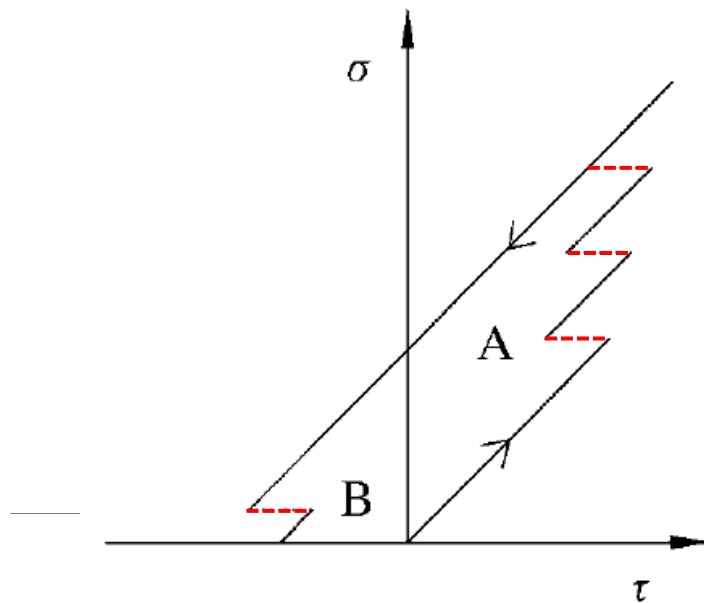
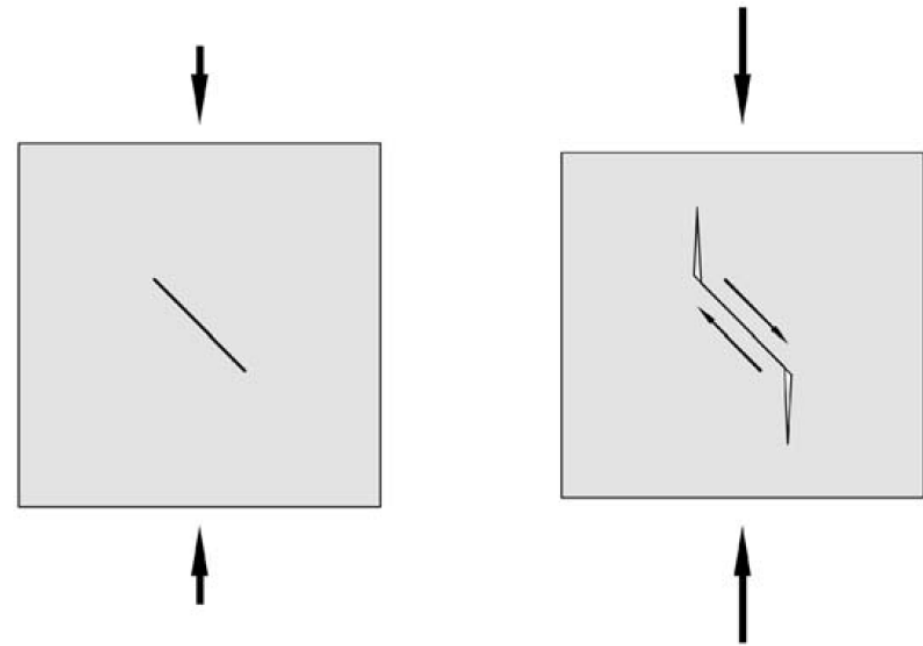
Cracks are sensitive to changes in stress

A compressive principal stress tends to close a crack that is oriented normal to the stress



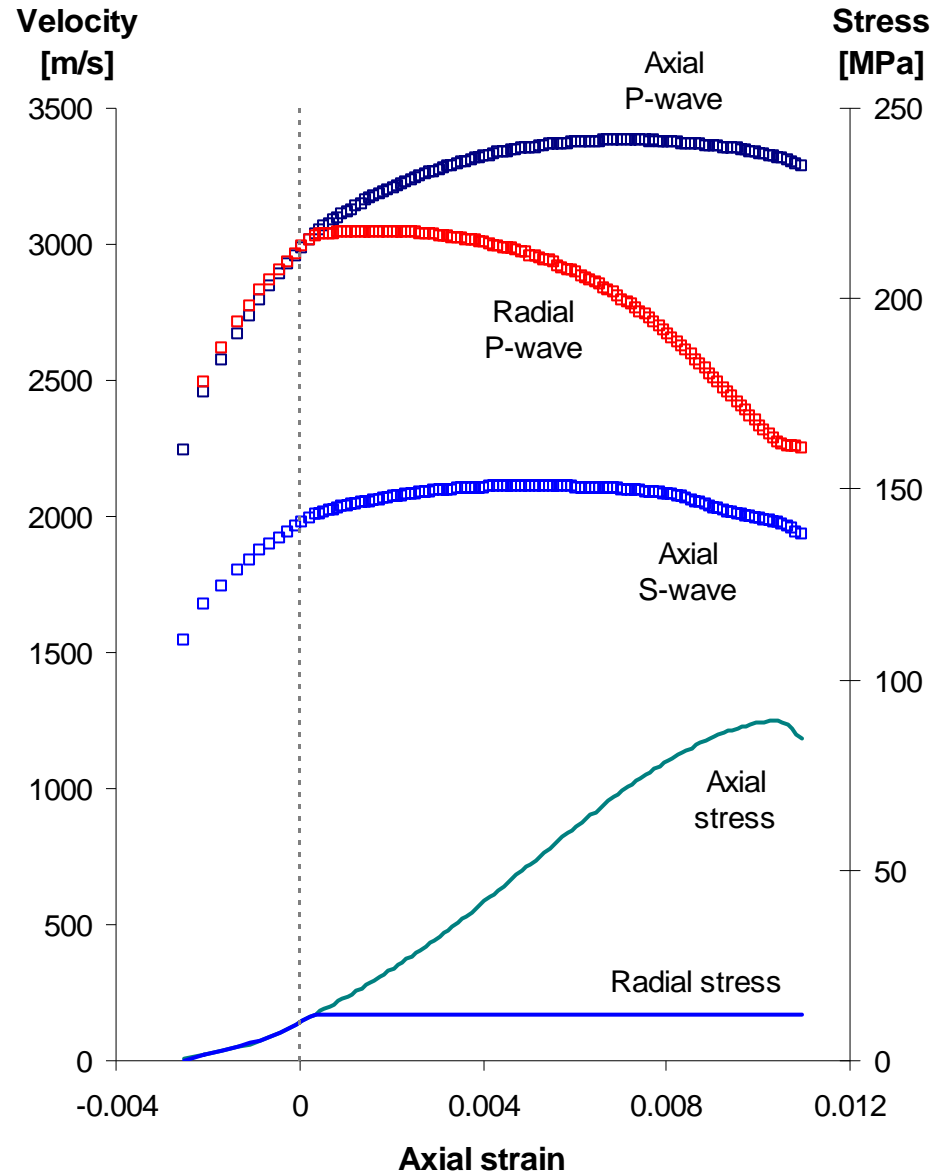
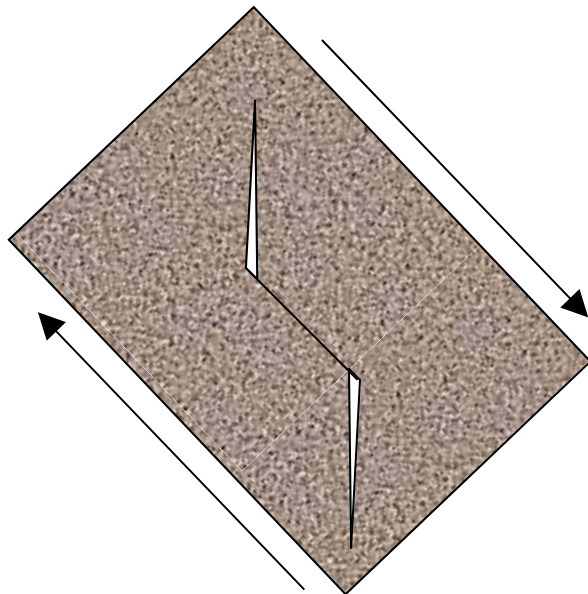
Sliding cracks

- induce hysteresis,
permanent deformation,
and difference between
loading and unloading
modulus



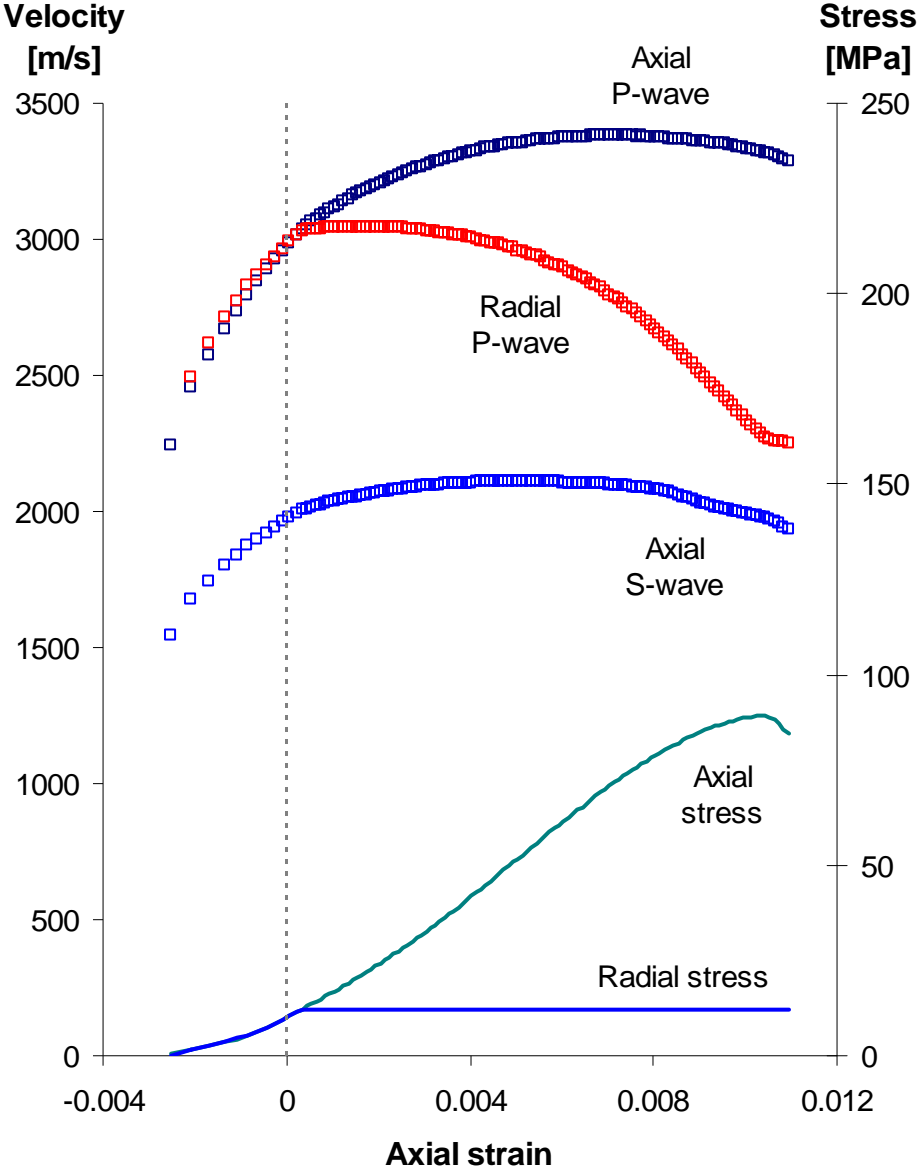
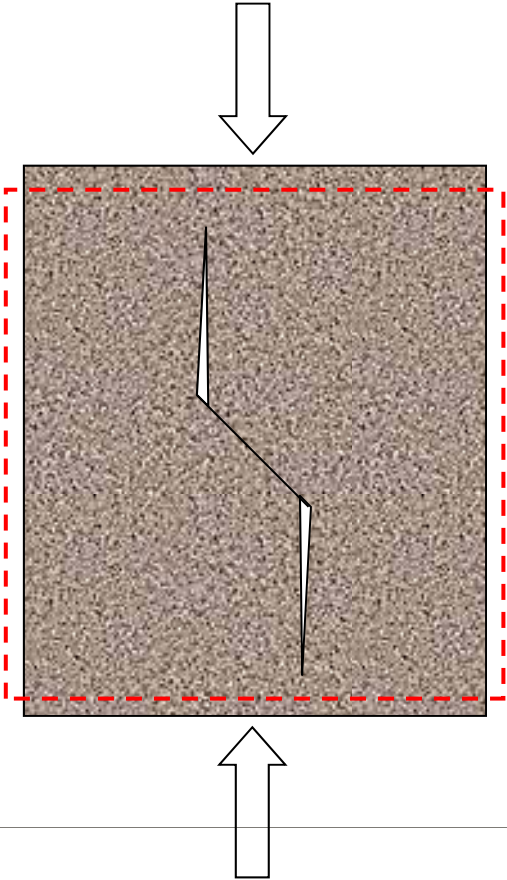
Cracks are sensitive to changes in stress

Shear deformations tend to open up cracks



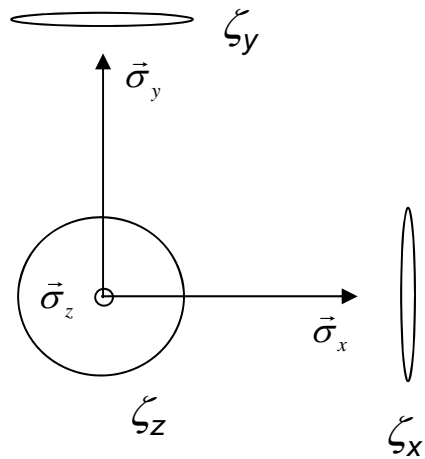
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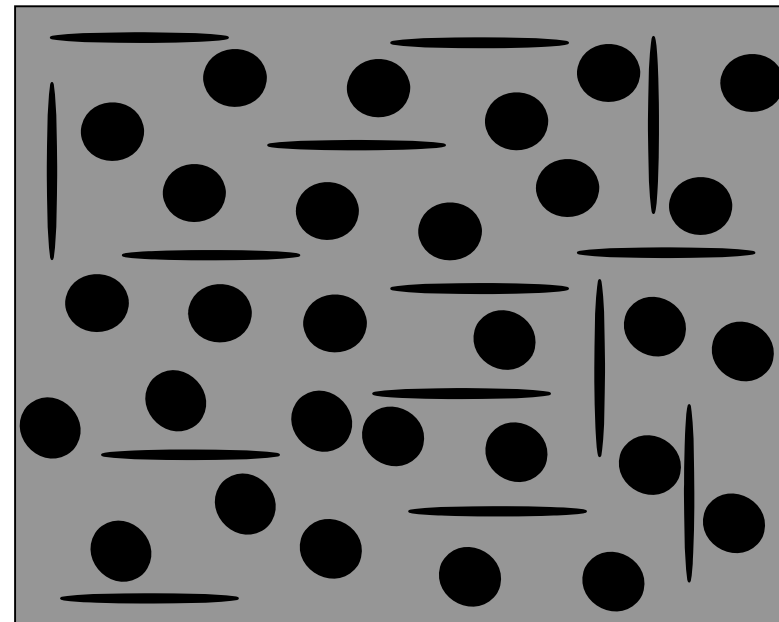


Fjær (2006):

Three sets of flat cracks
oriented normal to
the principal stresses



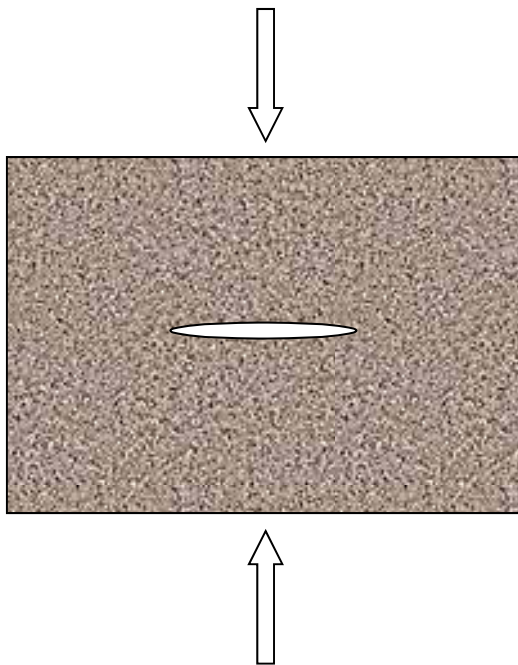
A model



Solid, pores & cracks

Cracks are sensitive to changes in stress

A compressive principal stress tends to close a crack that is oriented normal to the stress



We can not close a crack that is already closed

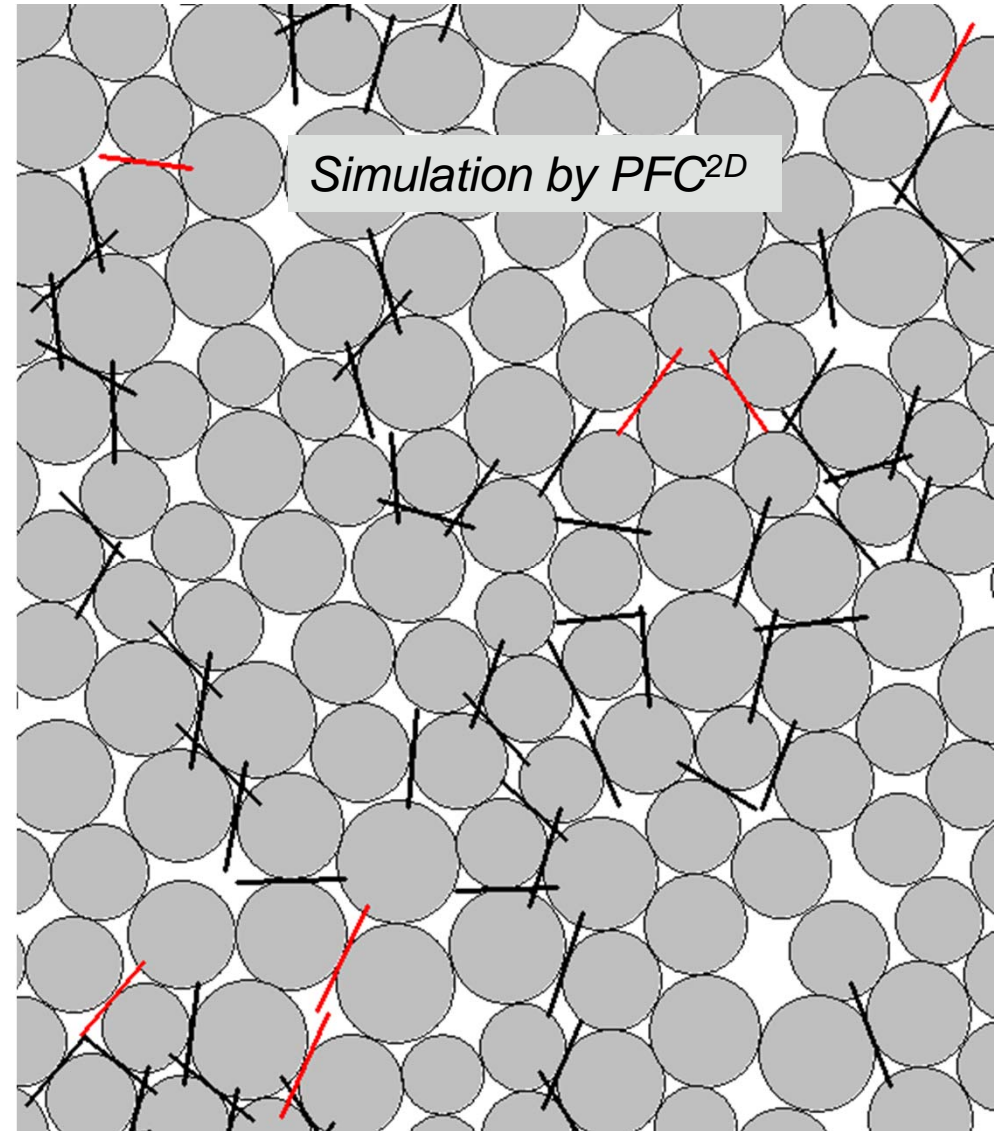
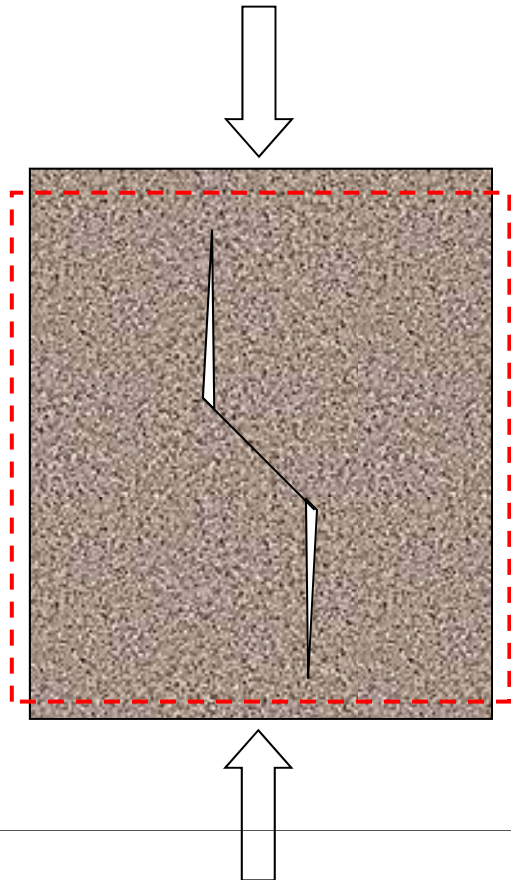
$$d\zeta_i \propto -\zeta_i d\sigma_i$$

Assumption:

$$\zeta_i \propto (\sigma_i + T_o)^{-n}$$

Cracks are sensitive to changes in stress

Shear deformations tend to open up cracks

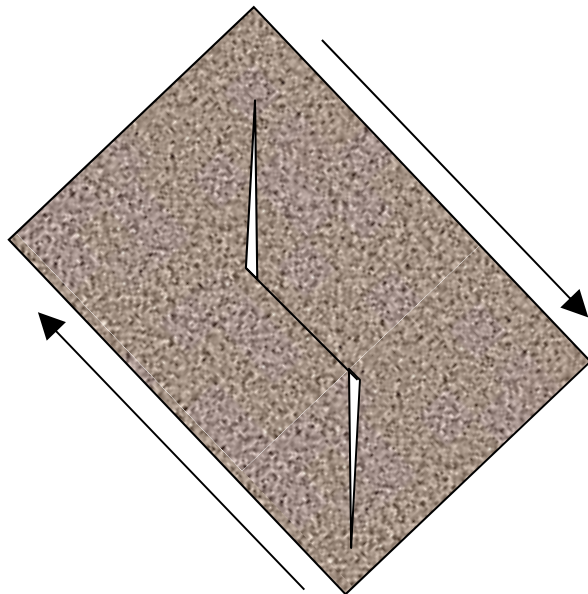


Local failures induced by compressive stress:

shear ———
tensile ———

Cracks are sensitive to changes in stress

Shear deformations tend to open up cracks



Sensitivity to shear strain:

$$d\zeta_i \propto -(d\varepsilon_i - d\varepsilon_j) - (d\varepsilon_i - d\varepsilon_k)$$

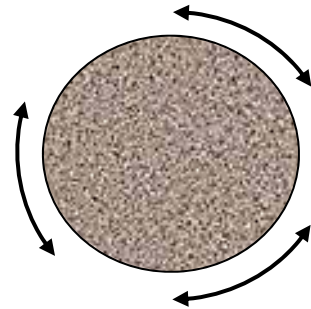
Assumption:

$$\zeta_i \propto e^{-\beta(2\varepsilon_i - \varepsilon_j - \varepsilon_k)}$$

Very large shear strains

⇒ more turbulent crack development

Changes in crack density more sensitive to magnitude than to orientation of shear strain



Assumption:

$$\zeta_i \propto e^{\Gamma^2}$$

Γ = maximum shear strain

Mathematics of the model:

$$C_{11} = C_{11}^o \left[1 - Q_{11}^p \phi - Q_{33} \zeta_x - Q_{11} (\zeta_y + \zeta_z) \right] \quad \text{etc.}$$

$$\zeta_i = \zeta_i^o \left(\frac{\sigma_i^o + T_o}{\sigma_i + T_o} \right)^n e^{-\beta(2\varepsilon_i - \varepsilon_j - \varepsilon_k) + \eta \Gamma^2}$$

$$\phi = \frac{\phi_o - \varepsilon_v}{1 - \varepsilon_v}$$

Velocities:

$$V_{p1} \left(= \sqrt{\frac{C_{11}}{\rho}} \right) \quad \text{etc.}$$

Assumptions:

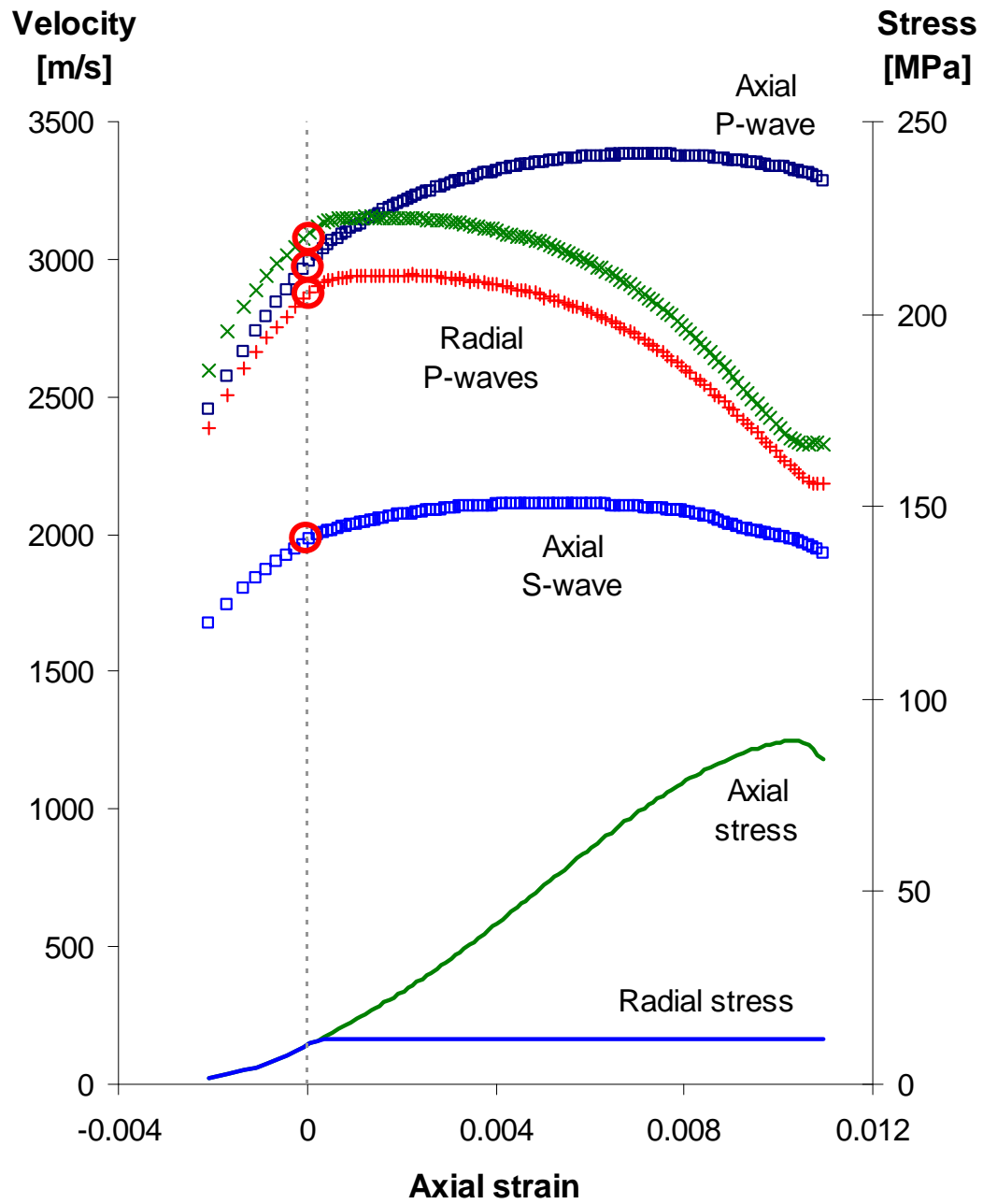
$$C_{11}^o = C_{22}^o = C_{33}^o = H_o$$

$$C_{44}^o = C_{55}^o = C_{66}^o = G_o$$

$$C_{12}^o = C_{13}^o = C_{23}^o = H_o - 2G_o$$

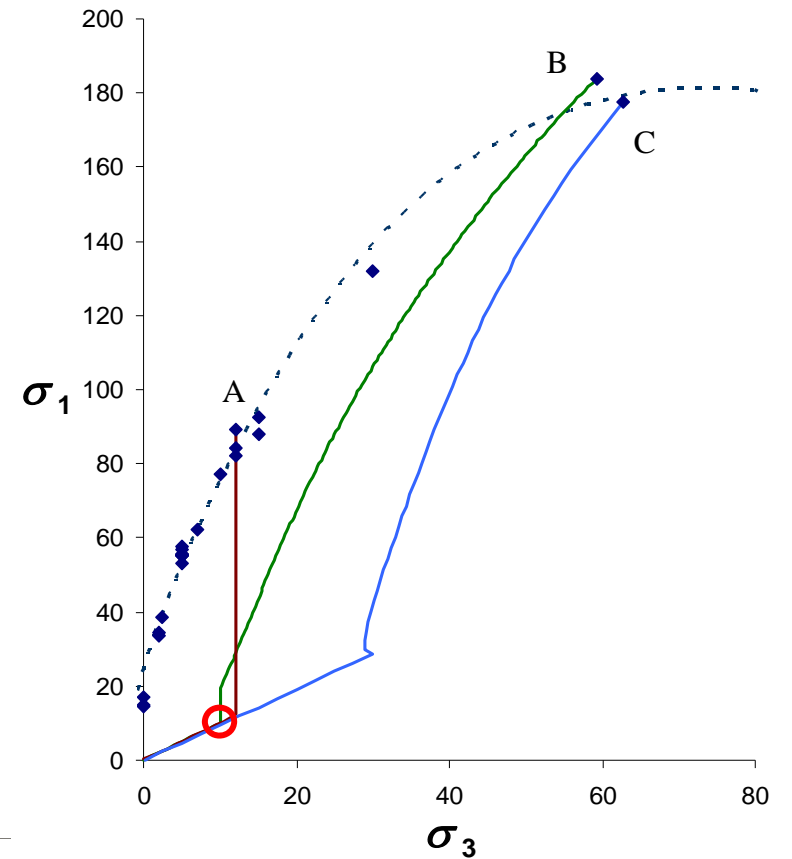
$$H_o = 80 \text{ GPa (fixed value)}$$

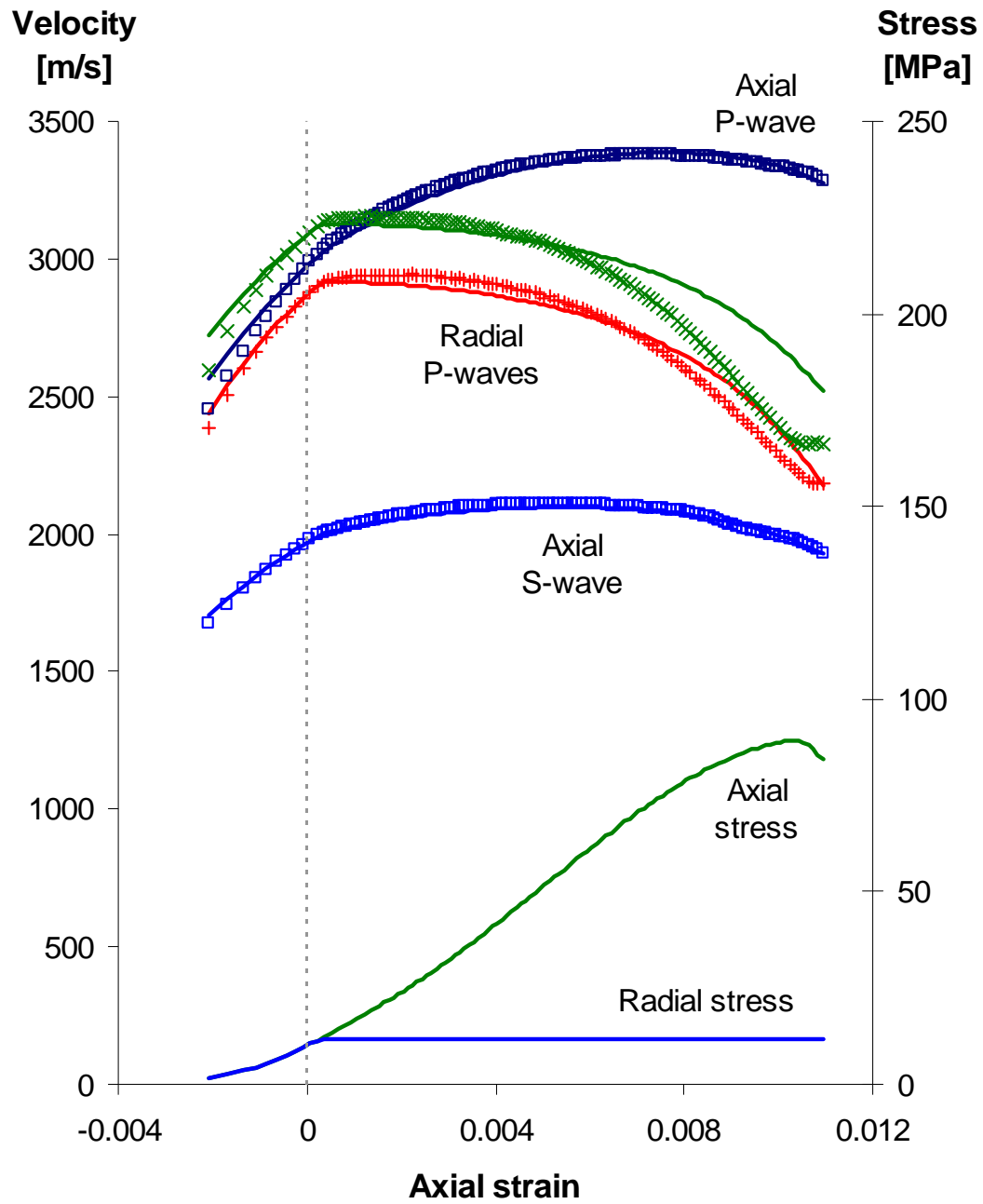
$$\nu = 0.2 \text{ (fixed value)}$$



Test A

○ Scaling: $\zeta_x^o, \zeta_y^o, \zeta_z^o, G_o$



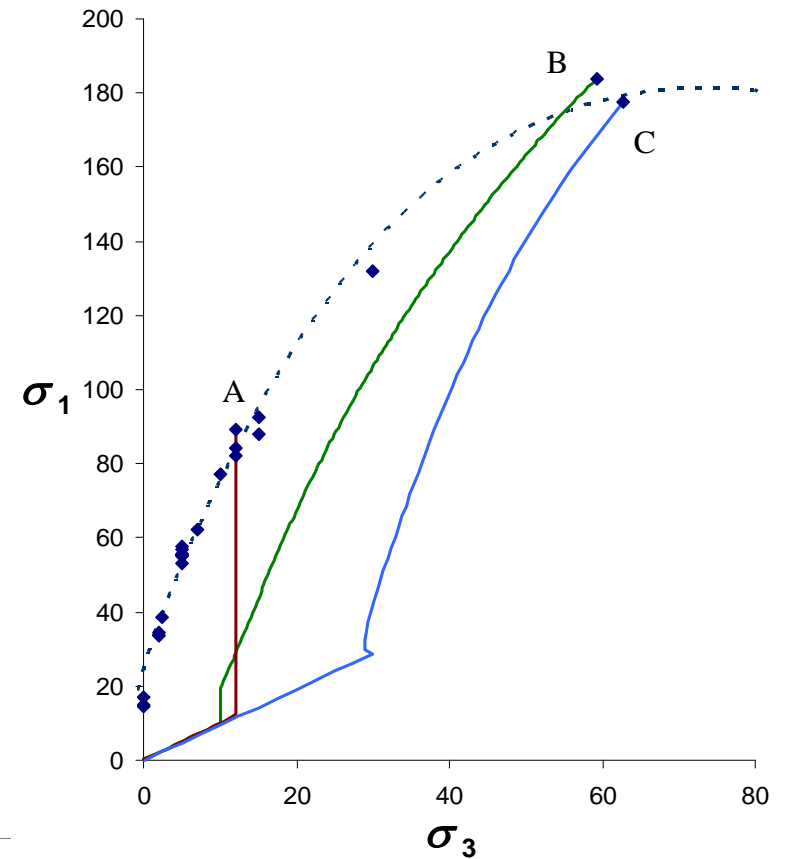


Test A

$$n = 0.1$$

$$\beta = 5$$

$$\eta = 800$$



⇒

The model – based on flat cracks and spherical pores - matches observations quite well

The match supports the claim that the stress dependency of wave velocities may largely be explained in terms of opening and closure of cracks

References:

Fjær, E., Holt, R.M., Horsrud, P., Raaen, A.M. and Risnes, R. (2008) "Petroleum Related Rock Mechanics. 2nd Edition". Elsevier, Amsterdam

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Sayers, C.M., Kachanov, M. (1995) "Microcrack-induced elastic wave anisotropy of brittle rocks", *J. Geophys. Res. B*, **100**, 4149-4156

Thomsen, L. (1995): "Elastic anisotropy due to aligned cracks in porous rock". *Geophysical Prospecting*, **43**, 805-829